

Stochastic differential equations with coefficients in Sobolev spaces

Shizan Fang^{c*}, Dejun Luo^{a,b}, Anton Thalmaier^a

^aUR Mathématiques, Université du Luxembourg, 6, rue Richard Coudenhove-Kalergi, L-1359 Luxembourg

^bKey Laboratory of Random Complex Structures and Data Science, Academy of Mathematics and Systems Science, Chinese Academy of Sciences, Beijing 100190, China

^cI.M.B, BP 47870, Université de Bourgogne, Dijon, France

Abstract

We consider Itô SDE $dX_t = \sum_{j=1}^m A_j(X_t) dw_t^j + A_0(X_t) dt$ on \mathbb{R}^d . The diffusion coefficients A_1, \dots, A_m are supposed to be in the Sobolev space $W_{\text{loc}}^{1,p}(\mathbb{R}^d)$ with $p > d$, and to have linear growth; for the drift coefficient A_0 , we consider two cases: (i) A_0 is continuous whose distributional divergence $\delta(A_0)$ w.r.t. the Gaussian measure γ_d exists, (ii) A_0 has the Sobolev regularity $W_{\text{loc}}^{1,p'}$ for some $p' > 1$. Assume $\int_{\mathbb{R}^d} \exp [\lambda_0 (|\delta(A_0)| + \sum_{j=1}^m (|\delta(A_j)|^2 + |\nabla A_j|^2))] d\gamma_d < +\infty$ for some $\lambda_0 > 0$, in the case (i), if the pathwise uniqueness of solutions holds, then the push-forward $(X_t)_{\#} \gamma_d$ admits a density with respect to γ_d . In particular, if the coefficients are bounded Lipschitz continuous, then X_t leaves the Lebesgue measure Leb_d quasi-invariant. In the case (ii), we develop a method used by G. Crippa and C. De Lellis for ODE and implemented by X. Zhang for SDE, to establish the existence and uniqueness of stochastic flow of maps.

MSC 2000: primary 60H10, 34F05; secondary 60J60, 37C10, 37H10.

Key words: Stochastic flows, Sobolev space coefficients, density, density estimate, pathwise uniqueness, Gaussian measure, Ornstein-Uhlenbeck semigroup.

1 Introduction

Let $A_0, A_1, \dots, A_m: \mathbb{R}^d \rightarrow \mathbb{R}^d$ be continuous vector fields on \mathbb{R}^d . We consider the following Itô stochastic differential equation on \mathbb{R}^d (abbreviated as SDE)

$$dX_t = \sum_{j=1}^m A_j(X_t) dw_t^j + A_0(X_t) dt, \quad X_0 = x, \quad (1.1)$$

where $w_t = (w_t^1, \dots, w_t^m)$ is the standard Brownian motion on \mathbb{R}^m . It is a classical fact in the theory of SDE (see [16, 17, 21, 30]) that, if the coefficients A_j are globally Lipschitz continuous, then SDE (1.1) has a unique strong solution which defines a stochastic flow of homeomorphisms on \mathbb{R}^d ; however contrary to ordinary differential equations (abbreviated as ODE), the regularity of the homeomorphisms is only Hölder continuity of order $0 < \alpha < 1$. Thus it is not clear whether the Lebesgue measure Leb_d on \mathbb{R}^d admits a density under the flow X_t . In the case where the vector fields $A_j, j = 0, 1, \dots, m$, are in $C_b^\infty(\mathbb{R}^d, \mathbb{R}^d)$, the SDE (1.1) defines a flow of diffeomorphisms, and Kunita [21] showed that the measures on \mathbb{R}^d which have a strictly positive

*fang@u-bourgogne.fr

smooth density with respect to Leb_d are quasi-invariant under the flow. This result was recently generalized in [27] to the case where the drift A_0 is allowed to be only log-Lipschitz continuous. Studies on SDE beyond the Lipschitz setting attracted great interest during the last years, see for instance [10, 11, 13, 19, 20, 23, 24, 29, 34, 35].

In the context of ODE, existence of a flow of quasi-invariant measurable maps associated to a vector field A_0 belonging to Sobolev spaces appeared first in [6]. In the seminar paper [7], Di Perna and Lions developed transport equations to solve ODE without involving exponential integrability of $|\nabla A_0|$. On the other hand, L. Ambrosio [1] took advantage of using continuity equations which allowed him to construct quasi-invariant flows associated to vector fields A_0 with only BV regularity. In the framework for Gaussian measures, the Di Perna-Lions method was developed in [4], also in [2, 12] on the Wiener space.

The situation for SDE is quite different: even for the vector fields A_0, A_1, \dots, A_m in C^∞ with linear growth, if no conditions were imposed on the growth of the derivatives, the SDE (1.1) could not define a flow of diffeomorphisms (see [25, 26]). More precisely, let τ_x be the life time of the solution to (1.1) starting from x . The SDE (1.1) is said to be *complete* if for each $x \in \mathbb{R}^d$, $\mathbb{P}(\tau_x = +\infty) = 1$; it is said to be *strongly complete* if $\mathbb{P}(\tau_x = +\infty, x \in \mathbb{R}^d) = 1$. The goal in [26] is to construct examples for which the coefficients are smooth, but the SDE (1.1) is not strongly complete (see [11, 25] for positive examples). Now consider

$$\Sigma = \{(w, x) \in \Omega \times \mathbb{R}^d; \tau_x(w) = +\infty\}.$$

Suppose that the SDE (1.1) is complete, then for any probability measure μ on \mathbb{R}^d ,

$$\int_{\mathbb{R}^d} \left(\int_{\Omega} \mathbf{1}_\Sigma(w, x) \, d\mathbb{P}(w) \right) d\mu(x) = 1.$$

By Fubini's theorem, $\int_{\Omega} \left(\int_{\mathbb{R}^d} \mathbf{1}_\Sigma(w, x) \, d\mu(x) \right) d\mathbb{P}(w) = 1$. It follows that there exists a full measure subset $\Omega_0 \subset \Omega$ such that for all $w \in \Omega_0$, $\tau_x(w) = +\infty$ holds for μ -almost every $x \in \mathbb{R}^d$. Now under the existence of a complete unique strong solution to SDE (1.1), we have a flow of measurable maps $x \rightarrow X_t(w, x)$.

Recently, inspired by a previous work due to Ambrosio, Lecumberry and Maniglia [3], Crippa and De Lellis [5] obtained some new type of estimates of perturbation for ODE whose coefficients have Sobolev regularity. More precisely, the absence of Lipschitz condition was filled by the following inequality: for $f \in W_{loc}^{1,1}(\mathbb{R}^d)$,

$$|f(x) - f(y)| \leq C_d |x - y| (M_R |\nabla f|(x) + M_R |\nabla f|(y))$$

holds for $x, y \in N^c$ and $|x - y| \leq R$, where N is a negligible set of \mathbb{R}^d and $M_R g$ is the maximal function defined by

$$M_R g(x) = \sup_{0 < r \leq R} \frac{1}{\text{Leb}_d(B(x, r))} \int_{B(x, r)} |g(y)| \, dy,$$

here $B(x, r) = \{y \in \mathbb{R}^d; |y - x| \leq r\}$; the classical moment estimate was replaced by estimating the quantity

$$\int_{B(0, r)} \log \left(\frac{|X_t(x) - \tilde{X}_t(x)|}{\sigma} + 1 \right) \, dx,$$

where $\sigma > 0$ is a small parameter. This method has recently been successfully implemented to SDE by X. Zhang in [36].

The aim in this paper is two-fold: first we shall study absolute continuity of the push-forward measure $(X_t)_\# \text{Leb}_d$ with respect to Leb_d , once the SDE (1.1) has a unique strong solution;

secondly we shall construct strong solutions (for almost all initial values) using the approach mentioned above for SDE with coefficients in Sobolev space. The key point is to obtain *a priori* L^p estimate for the density. To this end, we shall work with the standard Gaussian measure γ_d ; this will be done in Section 2. The main result in Section 3 is the following

Theorem 1.1. *Let A_0, A_1, \dots, A_m be continuous vector fields on \mathbb{R}^d of linear growth. Assume that the diffusion coefficients A_1, \dots, A_m are in the Sobolev space $\cap_{q>1} \mathbb{D}_1^q(\gamma_d)$ and that $\delta(A_0)$ exists; furthermore there exists a constant $\lambda_0 > 0$ such that*

$$\int_{\mathbb{R}^d} \exp \left[\lambda_0 \left(|\delta(A_0)| + \sum_{j=1}^m (|\delta(A_j)|^2 + |\nabla A_j|^2) \right) \right] d\gamma_d < +\infty. \quad (1.2)$$

Suppose that pathwise uniqueness holds for SDE (1.1). Then $(X_t)_{\#}\gamma_d$ is absolutely continuous with respect to γ_d and the density is in the space $L^1 \log L^1$.

A consequence of this theorem concerns the following classical situation.

Theorem 1.2. *Let A_0, A_1, \dots, A_m be globally Lipschitz continuous. Suppose that there exists a constant $C > 0$ such that*

$$\sum_{j=1}^m \langle x, A_j(x) \rangle^2 \leq C (1 + |x|^2) \quad \text{for all } x \in \mathbb{R}^d. \quad (1.3)$$

Then the stochastic flow of homeomorphisms X_t generated by SDE (1.1) leaves the Lebesgue measure Leb_d quasi-invariant.

Remark that the condition (1.3) not only includes the case of bounded Lipschitz diffusion coefficients, but also, maybe more significant, indicates the role of dispersion: the vector fields A_1, \dots, A_m should not go radically into infinity. The purpose of Section 4 is to find conditions that guarantee strict positivity of the density, in the case where the existence of the inverse flow is not known, see Theorem 4.4.

The main result in Section 5 is

Theorem 1.3. *Assume that the diffusion coefficients A_1, \dots, A_m belong to the Sobolev space $\cap_{q>1} \mathbb{D}_1^q(\gamma_d)$ and the drift $A_0 \in \mathbb{D}_1^q(\gamma_d)$ for some $q > 1$. Assume (1.2) and that the coefficients A_0, A_1, \dots, A_m are of linear growth, then there is a unique stochastic flow of measurable maps $X : [0, T] \times \Omega \times \mathbb{R}^d \rightarrow \mathbb{R}^d$, which solves (1.1) for almost all initial $x \in \mathbb{R}^d$ and the push-forward $(X_t(w, \cdot))_{\#}\gamma_d$ admits a density with respect to γ_d , which is in $L^1 \log L^1$.*

When the diffusion coefficients satisfy the uniform ellipticity, a classical result due to Stroock and Varadhan [32] says that if the diffusion coefficients A_1, \dots, A_m are bounded continuous and the drift A_0 is bounded Borel measurable, then the weak uniqueness holds, that is the uniqueness in law of the diffusion. This result was strengthened by Veretennikov [33], saying that in fact the pathwise uniqueness holds. When A_0 is not bounded, some conditions on diffusion coefficients were needed. In the case where the diffusion matrix $a = (a_{ij})$ is the identity, the drift A_0 in (1.1) can be quite singular: $A_0 \in L_{loc}^p(\mathbb{R}^d)$ with $p > d + 2$ implies that the SDE (1.1) has the pathwise uniqueness (see Krylov-Röckner [20] for a more complete study); if the diffusion coefficients A_1, \dots, A_m are bounded continuous, under a Sobolev condition, namely, $A_j \in W_{loc}^{1,2(d+1)}$ for $j = 1, \dots, m$ and $A_0 \in L_{loc}^{2(d+1)}(\mathbb{R}^d)$, X. Zhang proved in [34] that the SDE (1.1) admits a unique strong solution. Note that even in this uniformly non-degenerated case, if the diffusion coefficients lose the continuity, there are counterexamples for which the weak uniqueness does not hold, see [19, 31].

Finally we would like to mention that under weaker Sobolev type conditions, the connection between weak solutions and Fokker-Planck equations was investigated in [14, 22], some notions of “generalized solutions”, as well as the phenomena of coalescence and splitting, were investigated in [23, 24]. Stochastic transport equations were studied in [15, 36].

2 L^p estimate of the density

The purpose of this section is to derive *a priori* estimates for the density; we assume that the coefficients A_0, A_1, \dots, A_m of SDE (1.1) are *smooth with compact support* in \mathbb{R}^d . Then the solution X_t , i.e., $x \mapsto X_t(x)$, is a stochastic flow of diffeomorphisms on \mathbb{R}^d . Moreover SDE (1.1) is equivalent to the following Stratonovich SDE

$$dX_t = \sum_{j=1}^m A_j(X_t) \circ dw_t^j + \tilde{A}_0(X_t) dt, \quad X_0 = x, \quad (2.1)$$

where $\tilde{A}_0 = A_0 - \frac{1}{2} \sum_{j=1}^m \mathcal{L}_{A_j} A_j$ and \mathcal{L}_A denotes the Lie derivative with respect to A .

Let γ_d be the standard Gaussian measure on \mathbb{R}^d , and $\gamma_t = (X_t)_\# \gamma_d$, $\tilde{\gamma}_t = (X_t^{-1})_\# \gamma_d$ the push-forwards of γ_d respectively by the flow X_t and its inverse flow X_t^{-1} . To fix ideas, we denote by $(\Omega, \mathcal{F}, \mathbb{P})$ the probability space on which the Brownian motion w_t is defined. Let $K_t = \frac{d\gamma_t}{d\gamma_d}$ and $\tilde{K}_t = \frac{d\tilde{\gamma}_t}{d\gamma_d}$ be the densities with respect to γ_d . By Lemma 4.3.1 in [21], the Radon-Nikodym derivative \tilde{K}_t has the following explicit expression

$$\tilde{K}_t(x) = \exp \left(- \sum_{j=1}^m \int_0^t \delta(A_j)(X_s(x)) \circ dw_s^j - \int_0^t \delta(\tilde{A}_0)(X_s(x)) ds \right), \quad (2.2)$$

where $\delta(A_j)$ denotes the divergence of A_j with respect to the Gaussian measure γ_d :

$$\int_{\mathbb{R}^d} \langle \nabla \varphi, A_j \rangle d\gamma_d = \int_{\mathbb{R}^d} \varphi \delta(A_j) d\gamma_d, \quad \varphi \in C_c^1(\mathbb{R}^d).$$

It is easy to see that K_t and \tilde{K}_t are related to each other by the equality below:

$$K_t(x) = [\tilde{K}_t(X_t^{-1}(x))]^{-1}. \quad (2.3)$$

In fact, for any $\psi \in C_c^\infty(\mathbb{R}^d)$, we have

$$\begin{aligned} \int_{\mathbb{R}^d} \psi(x) d\gamma_d(x) &= \int_{\mathbb{R}^d} \psi[X_t(X_t^{-1}(x))] d\gamma_d(x) \\ &= \int_{\mathbb{R}^d} \psi[X_t(y)] \tilde{K}_t(y) d\gamma_d(y) = \int_{\mathbb{R}^d} \psi(x) \tilde{K}_t(X_t^{-1}(x)) K_t(x) d\gamma_d(x), \end{aligned}$$

which leads to (2.3) due to the arbitrariness of $\psi \in C_c^\infty(\mathbb{R}^d)$. In the following we shall estimate the $L^p(\mathbb{P} \times \gamma_d)$ norm of K_t .

We rewrite the density (2.2) with the Itô integral:

$$\tilde{K}_t(x) = \exp \left(- \sum_{j=1}^m \int_0^t \delta(A_j)(X_s(x)) dw_s^j - \int_0^t \left[\frac{1}{2} \sum_{j=1}^m \mathcal{L}_{A_j} \delta(A_j) + \delta(\tilde{A}_0) \right] (X_s(x)) ds \right). \quad (2.4)$$

Lemma 2.1. *We have*

$$\frac{1}{2} \sum_{j=1}^m \mathcal{L}_{A_j} \delta(A_j) + \delta(\tilde{A}_0) = \delta(A_0) + \frac{1}{2} \sum_{j=1}^m |A_j|^2 + \frac{1}{2} \sum_{j=1}^m \langle \nabla A_j, (\nabla A_j)^* \rangle, \quad (2.5)$$

where $\langle \cdot, \cdot \rangle$ denotes the inner product of $\mathbb{R}^d \otimes \mathbb{R}^d$ and $(\nabla A_j)^*$ the transpose of ∇A_j .

Proof. Let A be a C^2 vector field on \mathbb{R}^d . From the expression

$$\delta(A) = \sum_{k=1}^d \left(x_k A^k - \frac{\partial A^k}{\partial x_k} \right),$$

we get

$$\mathcal{L}_A \delta(A) = \sum_{\ell, k=1}^d \left(A^\ell A^k \delta_{k\ell} + A^\ell x_k \frac{\partial A^k}{\partial x_\ell} - A^\ell \frac{\partial^2 A^k}{\partial x_\ell \partial x_k} \right). \quad (2.6)$$

Note that

$$\frac{\partial}{\partial x_k} \left(A^\ell \frac{\partial A^k}{\partial x_\ell} \right) = \frac{\partial A^k}{\partial x_\ell} \frac{\partial A^\ell}{\partial x_k} + A^\ell \frac{\partial^2 A^k}{\partial x_k \partial x_\ell}.$$

Thus, by means of (2.6), we obtain

$$\mathcal{L}_A \delta(A) = |A|^2 + \delta(\mathcal{L}_A A) + \langle \nabla A, (\nabla A)^* \rangle. \quad (2.7)$$

Recall that $\delta(\tilde{A}_0) = \delta(A_0) - \frac{1}{2} \sum_{j=1}^m \delta(\mathcal{L}_{A_j} A_j)$. Hence, replacing A by A_j in (2.7) and summing over j , gives formula (2.5). \square

We can now prove the following key estimate.

Theorem 2.2. *For $p > 1$,*

$$\|K_t\|_{L^p(\mathbb{P} \times \gamma_d)} \leq \left[\int_{\mathbb{R}^d} \exp \left(pt \left[2|\delta(A_0)| + \sum_{j=1}^m (|A_j|^2 + |\nabla A_j|^2 + 2(p-1)|\delta(A_j)|^2) \right] \right) d\gamma_d \right]^{\frac{p-1}{p(2p-1)}}. \quad (2.8)$$

Proof. Using relation (2.3), we have

$$\begin{aligned} \int_{\mathbb{R}^d} \mathbb{E}[K_t^p(x)] d\gamma_d(x) &= \mathbb{E} \int_{\mathbb{R}^d} [\tilde{K}_t(X_t^{-1}(x))]^{-p} d\gamma_d(x) \\ &= \mathbb{E} \int_{\mathbb{R}^d} [\tilde{K}_t(y)]^{-p} \tilde{K}_t(y) d\gamma_d(y) \\ &= \int_{\mathbb{R}^d} \mathbb{E}[(\tilde{K}_t(x))^{-p+1}] d\gamma_d(x). \end{aligned} \quad (2.9)$$

To simplify the notation, denote the right hand side of (2.5) by Φ . Then $\tilde{K}_t(x)$ rewrites as

$$\tilde{K}_t(x) = \exp \left(- \sum_{j=1}^m \int_0^t \delta(A_j)(X_s(x)) dw_s^j - \int_0^t \Phi(X_s(x)) ds \right).$$

Fixing an arbitrary $r > 0$, we get

$$\begin{aligned}
(\tilde{K}_t(x))^{-r} &= \exp \left(r \sum_{j=1}^m \int_0^t \delta(A_j)(X_s(x)) \, dw_s^j + r \int_0^t \Phi(X_s(x)) \, ds \right) \\
&= \exp \left(r \sum_{j=1}^m \int_0^t \delta(A_j)(X_s(x)) \, dw_s^j - r^2 \sum_{j=1}^m \int_0^t |\delta(A_j)(X_s(x))|^2 \, ds \right) \\
&\quad \times \exp \left(\int_0^t \left(r^2 \sum_{j=1}^m |\delta(A_j)|^2 + r\Phi \right) (X_s(x)) \, ds \right).
\end{aligned}$$

By Cauchy-Schwarz's inequality,

$$\begin{aligned}
\mathbb{E}[(\tilde{K}_t(x))^{-r}] &\leq \left[\mathbb{E} \exp \left(2r \sum_{j=1}^m \int_0^t \delta(A_j)(X_s(x)) \, dw_s^j - 2r^2 \sum_{j=1}^m \int_0^t |\delta(A_j)(X_s(x))|^2 \, ds \right) \right]^{1/2} \\
&\quad \times \left[\mathbb{E} \exp \left(\int_0^t \left(2r^2 \sum_{j=1}^m |\delta(A_j)|^2 + 2r\Phi \right) (X_s(x)) \, ds \right) \right]^{1/2} \\
&= \left[\mathbb{E} \exp \left(\int_0^t \left(2r^2 \sum_{j=1}^m |\delta(A_j)|^2 + 2r\Phi \right) (X_s(x)) \, ds \right) \right]^{1/2}, \tag{2.10}
\end{aligned}$$

since the first term on the right hand side of the inequality in (2.10) is the expectation of a martingale. Let

$$\tilde{\Phi}_r = 2r|\delta(A_0)| + r \sum_{j=1}^m (|A_j|^2 + |\nabla A_j|^2 + 2r|\delta(A_j)|^2).$$

Then by (2.10), along with the definition of Φ and Cauchy-Schwarz's inequality, we obtain

$$\int_{\mathbb{R}^d} \mathbb{E}[(\tilde{K}_t(x))^{-r}] \, d\gamma_d \leq \left[\int_{\mathbb{R}^d} \mathbb{E} \exp \left(\int_0^t \tilde{\Phi}_r(X_s(x)) \, ds \right) \, d\gamma_d \right]^{1/2}. \tag{2.11}$$

Following the idea of A.B. Cruzeiro ([6] Corollary 2.2, see also Theorem 7.3 in [8]) and by Jensen's inequality,

$$\exp \left(\int_0^t \tilde{\Phi}_r(X_s(x)) \, ds \right) = \exp \left(\int_0^t t \tilde{\Phi}_r(X_s(x)) \frac{ds}{t} \right) \leq \frac{1}{t} \int_0^t e^{t \tilde{\Phi}_r(X_s(x))} \, ds.$$

Define $I(t) = \sup_{0 \leq s \leq t} \int_{\mathbb{R}^d} \mathbb{E}[K_t^p(x)] \, d\gamma_d$. Integrating on both sides of the above inequality and by Hölder's inequality,

$$\begin{aligned}
\int_{\mathbb{R}^d} \mathbb{E} \exp \left(\int_0^t \tilde{\Phi}_r(X_s(x)) \, ds \right) \, d\gamma_d(x) &\leq \frac{1}{t} \int_0^t \mathbb{E} \int_{\mathbb{R}^d} e^{t \tilde{\Phi}_r(X_s(x))} \, d\gamma_d(x) \, ds \\
&= \frac{1}{t} \int_0^t \mathbb{E} \int_{\mathbb{R}^d} e^{t \tilde{\Phi}_r(y)} K_s(y) \, d\gamma_d(y) \, ds \\
&\leq \frac{1}{t} \int_0^t \|e^{t \tilde{\Phi}_r}\|_{L^q(\gamma_d)} \|K_s\|_{L^p(\mathbb{P} \times \gamma_d)} \, ds \\
&\leq \|e^{t \tilde{\Phi}_r}\|_{L^q(\gamma_d)} I(t)^{1/p},
\end{aligned}$$

where q is the conjugate number of p . Thus it follows from (2.11) that

$$\int_{\mathbb{R}^d} \mathbb{E}[(\tilde{K}_t(x))^{-r}] d\gamma_d(x) \leq \|e^{t\tilde{\Phi}_r}\|_{L^q(\gamma_d)}^{1/2} I(t)^{1/2p}. \quad (2.12)$$

Taking $r = p - 1$ in the above estimate and by (2.9), we obtain

$$\int_{\mathbb{R}^d} \mathbb{E}[K_t^p(x)] d\gamma_d(x) \leq \|e^{t\tilde{\Phi}_{p-1}}\|_{L^q(\gamma_d)}^{1/2} I(t)^{1/2p}.$$

Thus we have $I(t) \leq \|e^{t\tilde{\Phi}_{p-1}}\|_{L^q(\gamma_d)}^{1/2} I(t)^{1/2p}$. Solving this inequality for $I(t)$ gives

$$\int_{\mathbb{R}^d} \mathbb{E}[K_t^p(x)] d\gamma_d(x) \leq I(t) \leq \left[\int_{\mathbb{R}^d} \exp\left(\frac{pt}{p-1}\tilde{\Phi}_{p-1}(x)\right) d\gamma_d(x) \right]^{\frac{p-1}{2p-1}}.$$

Now the desired estimate follows from the definition of $\tilde{\Phi}_{p-1}$. \square

Corollary 2.3. *For any $p > 1$,*

$$\|\tilde{K}_t\|_{L^p(\mathbb{P} \times \gamma_d)} \leq \left[\int_{\mathbb{R}^d} \exp\left((p+1)t\left[2|\delta(A_0)| + \sum_{j=1}^m (|A_j|^2 + |\nabla A_j|^2 + 2p|\delta(A_j)|^2)\right]\right) d\gamma_d \right]^{\frac{1}{2p+1}}. \quad (2.13)$$

Proof. Similar to (2.12), we have for $r > 0$,

$$\int_{\mathbb{R}^d} \mathbb{E}[(\tilde{K}_t(x))^r] d\gamma_d(x) \leq \|e^{t\tilde{\Phi}_r}\|_{L^q(\gamma_d)}^{1/2} I(t)^{1/2p}, \quad (2.14)$$

where $\tilde{\Phi}_r$ and $I(t)$ are defined as above. Since $I(t) \leq \|e^{t\tilde{\Phi}_{p-1}}\|_{L^q(\gamma_d)}^{p/(2p-1)}$, by taking $r = p - 1$, we get

$$\begin{aligned} \int_{\mathbb{R}^d} \mathbb{E}[(\tilde{K}_t(x))^{p-1}] d\gamma_d(x) &\leq \|e^{t\tilde{\Phi}_{p-1}}\|_{L^q(\gamma_d)}^{p/(2p-1)} \\ &= \left[\int_{\mathbb{R}^d} \exp\left(pt\left[2|\delta(A_0)| + \sum_{j=1}^m (|A_j|^2 + |\nabla A_j|^2 + 2(p-1)|\delta(A_j)|^2)\right]\right) d\gamma_d \right]^{\frac{p-1}{2p-1}}. \end{aligned}$$

Replacing p by $p + 1$ in the last inequality gives the claimed estimate. \square

3 Absolute continuity under flows generated by SDEs

Now assume that the coefficients A_j in SDE (1.1) are *continuous* and of linear growth. Then it is well known that SDE (1.1) has a weak solution of infinite life time. In order to apply the results of the preceding section, we shall regularize the vector fields using the Ornstein-Uhlenbeck semigroup $\{P_\varepsilon\}_{\varepsilon>0}$ on \mathbb{R}^d :

$$P_\varepsilon A(x) = \int_{\mathbb{R}^d} A(e^{-\varepsilon}x + \sqrt{1-e^{-2\varepsilon}}y) d\gamma_d(y).$$

We have the following simple properties.

Lemma 3.1. *Assume that A is continuous and $|A(x)| \leq C(1 + |x|^q)$ for some $q \geq 0$. Then*

(i) there is $C_q > 0$ independent of ε , such that

$$|P_\varepsilon A(x)| \leq C_q (1 + |x|^q), \quad \text{for all } x \in \mathbb{R}^d;$$

(ii) $P_\varepsilon A$ converges uniformly to A on any compact subset as $\varepsilon \rightarrow 0$.

Proof. (i) Note that $|e^{-\varepsilon}x + \sqrt{1 - e^{-2\varepsilon}}y| \leq |x| + |y|$ and that there exists a constant $C > 0$ such that $(|x| + |y|)^q \leq C(|x|^q + |y|^q)$. Using the growth condition on A , we have for some constant $C > 0$ (depending on q),

$$\begin{aligned} |P_\varepsilon A(x)| &\leq \int_{\mathbb{R}^d} |A(e^{-\varepsilon}x + \sqrt{1 - e^{-2\varepsilon}}y)| d\gamma_d(y) \\ &\leq C \int_{\mathbb{R}^d} (1 + |x|^q + |y|^q) d\gamma_d(y) \leq C(1 + |x|^q + M_q) \end{aligned}$$

where $M_q = \int_{\mathbb{R}^d} |y|^q d\gamma_d(y)$. Changing the constant yields (i).

(ii) Fix $R > 0$ and x in the closed ball $B(R)$ of radius R , centered at 0. Let $R_1 > R$ be arbitrary. We have

$$\begin{aligned} |P_\varepsilon A(x) - A(x)| &\leq \int_{\mathbb{R}^d} |A(e^{-\varepsilon}x + \sqrt{1 - e^{-2\varepsilon}}y) - A(x)| d\gamma_d(y) \\ &= \left(\int_{B(R_1)} + \int_{B(R_1)^c} \right) |A(e^{-\varepsilon}x + \sqrt{1 - e^{-2\varepsilon}}y) - A(x)| d\gamma_d(y) \\ &=: I_1 + I_2. \end{aligned} \tag{3.1}$$

By the growth condition on A , for some constant $C_q > 0$, independent of ε , we have

$$\begin{aligned} I_2 &\leq \int_{B(R_1)^c} \left(|A(e^{-\varepsilon}x + \sqrt{1 - e^{-2\varepsilon}}y)| + |A(x)| \right) d\gamma_d(y) \\ &\leq C_q \int_{B(R_1)^c} (1 + R^q + |y|^q) d\gamma_d(y), \end{aligned}$$

where the last term tends to 0 as $R_1 \rightarrow +\infty$. For given $\eta > 0$, we may take R_1 large enough such that $I_2 < \eta$. Then there exists $\varepsilon_{R_1} > 0$ such that for $\varepsilon < \varepsilon_{R_1}$ and $|y| \leq R_1$,

$$|e^{-\varepsilon}x + \sqrt{1 - e^{-2\varepsilon}}y| \leq e^{-\varepsilon}R + \sqrt{1 - e^{-2\varepsilon}}R_1 \leq R_1.$$

Note that

$$|e^{-\varepsilon}x + \sqrt{1 - e^{-2\varepsilon}}y - x| \leq \varepsilon R + \sqrt{2\varepsilon}R_1, \quad \text{for } |x| \leq R, |y| \leq R_1.$$

Since A is uniformly continuous on $B(R_1)$, there exists $\varepsilon_0 \leq \varepsilon_{R_1}$ such that

$$|A(e^{-\varepsilon}x + \sqrt{1 - e^{-2\varepsilon}}y) - A(x)| \leq \eta \quad \text{for all } y \in B(R_1), \varepsilon \leq \varepsilon_0.$$

As a result, the term $I_1 \leq \eta$. Therefore by (3.1), for any $\varepsilon \leq \varepsilon_0$,

$$\sup_{|x| \leq R} |P_\varepsilon A(x) - A(x)| \leq 2\eta.$$

The result follows from the arbitrariness of $\eta > 0$. □

The vector field $P_\varepsilon A$ is smooth on \mathbb{R}^d but does not have compact support. We introduce cut-off functions $\varphi_\varepsilon \in C_c^\infty(\mathbb{R}^d, [0, 1])$ satisfying

$$\varphi_\varepsilon(x) = 1 \quad \text{if } |x| \leq \frac{1}{\varepsilon}, \quad \varphi_\varepsilon(x) = 0 \quad \text{if } |x| \geq \frac{1}{\varepsilon} + 2 \quad \text{and } \|\nabla \varphi_\varepsilon\|_\infty \leq 1.$$

Set

$$A_j^\varepsilon = \varphi_\varepsilon P_\varepsilon A_j, \quad j = 0, 1, \dots, m.$$

Now consider the Itô SDE (1.1) with A_j being replaced by A_j^ε ($j = 0, 1, \dots, m$), and denote the corresponding terms by adding the superscript ε , e.g. X_t^ε , K_t^ε , etc.

In the sequel, we shall give a uniform estimate to K_t^ε . To this end, we need some preparations in the spirit of Malliavin calculus [28]. For a vector field A on \mathbb{R}^d and $p > 1$, we say that $A \in \mathbb{D}_1^p(\gamma_d)$ if $A \in L^p(\gamma_d)$ and if there exists $\nabla A: \mathbb{R}^d \rightarrow \mathbb{R}^d \otimes \mathbb{R}^d$ in $L^p(\gamma_d)$ such that for any $v \in \mathbb{R}^d$,

$$\nabla A(x)(v) = \partial_v A := \lim_{\eta \rightarrow 0} \frac{A(x + \eta v) - A(x)}{\eta} \quad \text{holds in } L^{p'}(\gamma_d) \text{ for any } p' < p.$$

For such $A \in \mathbb{D}_1^p(\gamma_d)$, the divergence $\delta(A) \in L^p(\gamma_d)$ exists and the following relations hold:

$$\nabla P_\varepsilon A = e^{-\varepsilon} P_\varepsilon (\nabla A), \quad \delta(P_\varepsilon A) = e^\varepsilon P_\varepsilon (\delta(A)). \quad (3.2)$$

If $A \in L^p(\gamma_d)$, then $P_\varepsilon A \in \mathbb{D}_1^p(\gamma_d)$ and $\lim_{\varepsilon \rightarrow 0} \|P_\varepsilon A - A\|_{L^p} = 0$.

Lemma 3.2. *Assume the vector field $A \in \mathbb{D}_1^p(\gamma_d)$ with $p > 1$, and denote by $A^\varepsilon = \varphi_\varepsilon P_\varepsilon A$. Then for $\varepsilon \in]0, 1]$,*

$$\begin{aligned} |\delta(A^\varepsilon)| &\leq P_\varepsilon(|A| + e|\delta(A)|), \\ |A^\varepsilon|^2 &\leq P_\varepsilon(|A|^2), \\ |\nabla A^\varepsilon|^2 &\leq P_\varepsilon[2(|A|^2 + |\nabla A|^2)], \\ |\delta(A^\varepsilon)|^2 &\leq P_\varepsilon[2(|A|^2 + e^2|\delta(A)|^2)]. \end{aligned}$$

Proof. Note that according to (3.2), $\delta(A^\varepsilon) = \delta(\varphi_\varepsilon P_\varepsilon A) = \varphi_\varepsilon e^\varepsilon P_\varepsilon \delta(A) - \langle \nabla \varphi_\varepsilon, P_\varepsilon A \rangle$, from where the first inequality follows. In the same way, the other results are obtained. \square

Applying Theorem 2.2 to K_t^ε with $p = 2$, we have

$$\|K_t^\varepsilon\|_{L^2(\mathbb{P} \times \gamma_d)} \leq \left[\int_{\mathbb{R}^d} \exp \left(2t \left[2|\delta(A_0^\varepsilon)| + \sum_{j=1}^m (|A_j^\varepsilon|^2 + |\nabla A_j^\varepsilon|^2 + 2|\delta(A_j^\varepsilon)|^2) \right] \right) d\gamma_d \right]^{1/6}. \quad (3.3)$$

By Lemma 3.2,

$$\begin{aligned} 2|\delta(A_0^\varepsilon)| + \sum_{j=1}^m (|A_j^\varepsilon|^2 + |\nabla A_j^\varepsilon|^2 + 2|\delta(A_j^\varepsilon)|^2) \\ \leq P_\varepsilon \left[2|A_0| + 2e|\delta(A_0)| + \sum_{j=1}^m (7|A_j|^2 + 2|\nabla A_j|^2 + 4e^2|\delta(A_j)|^2) \right]. \end{aligned}$$

We deduce from Jensen's inequality and the invariance of γ_d under the action of the semigroup P_ε that

$$\|K_t^\varepsilon\|_{L^2(\mathbb{P} \times \gamma_d)} \leq \left[\int_{\mathbb{R}^d} \exp \left(4t \left[|A_0| + e|\delta(A_0)| + \sum_{j=1}^m (4|A_j|^2 + |\nabla A_j|^2 + 2e^2|\delta(A_j)|^2) \right] \right) d\gamma_d \right]^{1/6} \quad (3.4)$$

for any $\varepsilon \leq 1$. According to (3.4), we consider the following conditions.

Assumptions (H):

(A1) For $j = 1, \dots, m$, $A_j \in \cap_{q \geq 1} \mathbb{D}_1^q(\gamma_d)$, A_0 is continuous and $\delta(A_0)$ exists.

(A2) The vector fields A_0, A_1, \dots, A_m have linear growth.

(A3) There exists $\lambda_0 > 0$ such that

$$\int_{\mathbb{R}^d} \exp \left[\lambda_0 \left(|\delta(A_0)| + \sum_{j=1}^m |\delta(A_j)|^2 \right) \right] d\gamma_d < +\infty.$$

(A4) There exists $\lambda_0 > 0$ such that

$$\int_{\mathbb{R}^d} \exp \left(\lambda_0 \sum_{j=1}^m |\nabla A_j|^2 \right) d\gamma_d < +\infty.$$

Note that by Sobolev's embedding theorem, the diffusion coefficients A_1, \dots, A_m admit Hölder continuous versions. In what follows, we consider these continuous versions. It is clear that under the conditions (A2)–(A4), there exists $T_0 > 0$ small enough, such that

$$\Lambda_{T_0} := \left[\int_{\mathbb{R}^d} \exp \left(4T_0 \left[|A_0| + e|\delta(A_0)| + \sum_{j=1}^m (4|A_j|^2 + |\nabla A_j|^2 + 2e^2|\delta(A_j)|^2) \right] \right) d\gamma_d \right]^{1/6} < \infty. \quad (3.5)$$

In this case, for $t \in [0, T_0]$,

$$\sup_{0 < \varepsilon \leq 1} \|K_t^\varepsilon\|_{L^2(\mathbb{P} \times \gamma_d)} \leq \Lambda_{T_0}. \quad (3.6)$$

Theorem 3.3. *Let $T > 0$ be given. Under (A1)–(A4) in Assumptions (H), there are two positive constants C_1 and C_2 , independent of ε , such that*

$$\sup_{0 < \varepsilon \leq 1} \mathbb{E} \int_{\mathbb{R}^d} K_t^\varepsilon |\log K_t^\varepsilon| d\gamma_d \leq 2(C_1 T)^{1/2} \Lambda_{T_0} + C_2 T \Lambda_{T_0}^2, \quad \text{for all } t \in [0, T].$$

Proof. We follow the arguments of Proposition 4.4 in [12]. By (2.3) and (2.4), we have

$$K_t^\varepsilon(X_t^\varepsilon(x)) = [\tilde{K}_t^\varepsilon(x)]^{-1} = \exp \left(\sum_{j=1}^m \int_0^t \delta(A_j^\varepsilon)(X_s^\varepsilon(x)) dw_s^j + \int_0^t \Phi_\varepsilon(X_s^\varepsilon(x)) ds \right),$$

where

$$\Phi_\varepsilon = \delta(A_0^\varepsilon) + \frac{1}{2} \sum_{j=1}^m |A_j^\varepsilon|^2 + \frac{1}{2} \sum_{j=1}^m \langle \nabla A_j^\varepsilon, (\nabla A_j^\varepsilon)^* \rangle.$$

Thus

$$\begin{aligned} \mathbb{E} \int_{\mathbb{R}^d} K_t^\varepsilon |\log K_t^\varepsilon| d\gamma_d &= \mathbb{E} \int_{\mathbb{R}^d} |\log K_t^\varepsilon(X_t^\varepsilon(x))| d\gamma_d(x) \\ &\leq \mathbb{E} \int_{\mathbb{R}^d} \left| \sum_{j=1}^m \int_0^t \delta(A_j^\varepsilon)(X_s^\varepsilon(x)) dw_s^j \right| d\gamma_d(x) + \mathbb{E} \int_{\mathbb{R}^d} \left| \int_0^t \Phi_\varepsilon(X_s^\varepsilon(x)) ds \right| d\gamma_d(x) \\ &=: I_1 + I_2. \end{aligned} \quad (3.7)$$

Using Burkholder's inequality, we get

$$\mathbb{E} \left| \sum_{j=1}^m \int_0^t \delta(A_j^\varepsilon)(X_s^\varepsilon(x)) dw_s^j \right| \leq 2 \mathbb{E} \left[\left(\sum_{j=1}^m \int_0^t |\delta(A_j^\varepsilon)(X_s^\varepsilon(x))|^2 ds \right)^{1/2} \right].$$

For the sake of simplifying the notations, write $\Psi_\varepsilon = \sum_{j=1}^m |\delta(A_j^\varepsilon)|^2$. By Cauchy's inequality,

$$I_1 \leq 2 \left[\int_0^t \mathbb{E} \int_{\mathbb{R}^d} |\Psi_\varepsilon(X_s^\varepsilon(x))| d\gamma_d(x) ds \right]^{1/2}. \quad (3.8)$$

Now we are going to estimate $\mathbb{E} \int_{\mathbb{R}^d} |\Psi_\varepsilon(X_s^\varepsilon(x))|^{2^\alpha} d\gamma_d(x)$ for $\alpha \in \mathbb{Z}_+$ which will be done inductively. First if $s \in [0, T_0]$, then by (3.4) and (3.6), along with Cauchy's inequality,

$$\begin{aligned} \mathbb{E} \int_{\mathbb{R}^d} |\Psi_\varepsilon(X_s^\varepsilon(x))|^{2^\alpha} d\gamma_d(x) &= \mathbb{E} \int_{\mathbb{R}^d} |\Psi_\varepsilon(y)|^{2^\alpha} K_s^\varepsilon(y) d\gamma_d(y) \\ &\leq \|\Psi_\varepsilon\|_{L^{2^\alpha+1}(\gamma_d)}^{2^\alpha} \|K_s^\varepsilon\|_{L^2(\mathbb{P} \times \gamma_d)} \\ &\leq \Lambda_{T_0} \|\Psi_\varepsilon\|_{L^{2^\alpha+1}(\gamma_d)}^{2^\alpha}. \end{aligned} \quad (3.9)$$

Now for $s \in [T_0, 2T_0]$, we shall use the flow property of X_s^ε : let $(\theta_{T_0} w)_t := w_{T_0+t} - w_{T_0}$ and X_t^{ε, T_0} be the solution of the Itô SDE driven by the new Brownian motion $(\theta_{T_0} w)_t$, then

$$X_{T_0+t}^\varepsilon(x, w) = X_t^{\varepsilon, T_0}(X_{T_0}^\varepsilon(x, w), \theta_{T_0} w), \quad \text{for all } t \geq 0,$$

and X_t^{ε, T_0} enjoys the same properties as X_t^ε . Therefore,

$$\begin{aligned} \mathbb{E} \int_{\mathbb{R}^d} |\Psi_\varepsilon(X_s^\varepsilon(x))|^{2^\alpha} d\gamma_d(x) &= \mathbb{E} \int_{\mathbb{R}^d} |\Psi_\varepsilon(X_{s-T_0}^{\varepsilon, T_0}(X_{T_0}^\varepsilon(x)))|^{2^\alpha} d\gamma_d(x) \\ &= \mathbb{E} \int_{\mathbb{R}^d} |\Psi_\varepsilon(X_{s-T_0}^{\varepsilon, T_0}(y))|^{2^\alpha} K_{T_0}^\varepsilon(y) d\gamma_d(y) \end{aligned}$$

which is dominated, using Cauchy-Schwarz inequality

$$\begin{aligned} &\left(\mathbb{E} \int_{\mathbb{R}^d} |\Psi_\varepsilon(X_{s-T_0}^{\varepsilon, T_0}(y))|^{2^{\alpha+1}} d\gamma_d(y) \right)^{1/2} \|K_{T_0}^\varepsilon\|_{L^2(\mathbb{P} \times \gamma_d)} \\ &\leq \left(\Lambda_{T_0} \|\Psi_\varepsilon\|_{L^{2^{\alpha+2}}(\gamma_d)}^{2^{\alpha+1}} \right)^{1/2} \Lambda_{T_0} = \Lambda_{T_0}^{1+2^{-1}} \|\Psi_\varepsilon\|_{L^{2^{\alpha+2}}(\gamma_d)}^{2^\alpha}. \end{aligned}$$

Repeating this procedure, we finally obtain, for all $s \in [0, T]$,

$$\mathbb{E} \int_{\mathbb{R}^d} |\Psi_\varepsilon(X_s^\varepsilon(x))|^{2^\alpha} d\gamma_d(x) \leq \Lambda_{T_0}^{1+2^{-1}+\dots+2^{-N+1}} \|\Psi_\varepsilon\|_{L^{2^{\alpha+N}}(\gamma_d)}^{2^\alpha},$$

where $N \in \mathbb{Z}_+$ is the unique integer such that $(N-1)T_0 < T \leq NT_0$. In particular, taking $\alpha = 0$ gives

$$\mathbb{E} \int_{\mathbb{R}^d} |\Psi_\varepsilon(X_s^\varepsilon(x))| d\gamma_d(x) \leq \Lambda_{T_0}^2 \|\Psi_\varepsilon\|_{L^{2^N}(\gamma_d)}. \quad (3.10)$$

By Lemma 3.2,

$$\sup_{0 < \varepsilon \leq 1} \|\Psi_\varepsilon\|_{L^{2^N}(\gamma_d)} \leq \left\| 2 \sum_{j=1}^m (|A_j|^2 + e^2 |\delta(A_j)|^2) \right\|_{L^{2^N}(\gamma_d)} =: C_1$$

whose right hand side is finite under the assumptions (A2)–(A4). This along with (3.8) and (3.10) leads to

$$I_1 \leq 2(C_1 T)^{1/2} \Lambda_{T_0}. \quad (3.11)$$

The same manipulation works for the term I_2 and we get

$$I_2 \leq C_2 T \Lambda_{T_0}^2, \quad (3.12)$$

where

$$C_2 = \left\| |A_0| + e|\delta(A_0)| + \frac{3}{2} \sum_{j=1}^m |A_j|^2 + \sum_{j=1}^m |\nabla A_j|^2 \right\|_{L^{2N}(\gamma_d)} < \infty.$$

Now we draw the conclusion from (3.7), (3.11) and (3.12). \square

It follows from Theorem 3.3 that the family $\{K^\varepsilon\}_{0 < \varepsilon \leq 1}$ is weakly compact in $L^1([0, T] \times \Omega \times \mathbb{R}^d)$. Along a subsequence, K^ε converges weakly to some $K \in L^1([0, T] \times \Omega \times \mathbb{R}^d)$ as $\varepsilon \rightarrow 0$. Let

$$\mathcal{C} = \left\{ u \in L^1([0, T] \times \Omega \times \mathbb{R}^d) : u_t \geq 0, \int_{\mathbb{R}^d} \mathbb{E}(u_t \log u_t) d\gamma_d \leq 2(C_1 T)^{1/2} \Lambda_{T_0} + C_2 T \Lambda_{T_0}^2 \right\}.$$

By convexity of the function $s \rightarrow s \log s$, it is clear that \mathcal{C} is a convex subset of $L^1([0, T] \times \Omega \times \mathbb{R}^d)$. Since the weak closure of \mathcal{C} coincides with the strong one, there exists a sequence of functions $u^{(n)} \in \mathcal{C}$ which converges to K in $L^1([0, T] \times \Omega \times \mathbb{R}^d)$. Along a subsequence, $u^{(n)}$ converges to K almost everywhere. Hence by Fatou's lemma, we get for almost all $t \in [0, T]$,

$$\int_{\mathbb{R}^d} \mathbb{E}(K_t \log K_t) d\gamma_d \leq 2(C_1 T)^{1/2} \Lambda_{T_0} + C_2 T \Lambda_{T_0}^2. \quad (3.13)$$

Theorem 3.4. *Assume conditions (A1)–(A4) and that pathwise uniqueness holds for SDE (1.1). Then for each $t > 0$, there is a full subset $\Omega_t \subset \Omega$ such that for all $w \in \Omega_t$, the density \hat{K}_t of $(X_t)_\# \gamma_d$ with respect to γ_d exists and $\hat{K}_t \in L^1 \log L^1$.*

Proof. Under these assumptions, we can use Theorem A in [18]. For the convenience of the reader, we include the statement:

Theorem 3.5 ([18]). *Let $\sigma_n(x)$ and $b_n(x)$ be continuous, taking values respectively in the space of $(d \times m)$ -matrices and \mathbb{R}^d . Suppose that*

$$\sup_n (\|\sigma_n(x)\| + |b_n(x)|) \leq C(1 + |x|),$$

and for any $R > 0$,

$$\lim_{n \rightarrow +\infty} \sup_{|x| \leq R} (\|\sigma_n(x) - \sigma(x)\| + |b_n(x) - b(x)|) = 0.$$

Suppose further that for the same Brownian motion $B(t)$, $X_n(x, t)$ solves the SDE

$$dX_n(t) = \sigma_n(X_n(t)) dB(t) + b_n(X_n(t)) dt, \quad X_n(0) = x.$$

If pathwise uniqueness holds for

$$dX(t) = \sigma(X(t)) dB(t) + b(X(t)) dt, \quad X(0) = x,$$

then for any $R > 0$, $T > 0$,

$$\lim_{n \rightarrow +\infty} \sup_{|x| \leq R} \mathbb{E} \left(\sup_{0 \leq t \leq T} |X_n(t, x) - X(t, x)|^2 \right) = 0. \quad (3.14)$$

We continue the proof of Theorem 3.4. By means of Lemma 3.1 and Theorem 3.5, for any $T, R > 0$, we get

$$\lim_{\varepsilon \rightarrow 0} \sup_{|x| \leq R} \mathbb{E} \left(\sup_{0 \leq t \leq T} |X_t^\varepsilon(x) - X_t(x)|^2 \right) = 0. \quad (3.15)$$

Now fixing arbitrary $\xi \in L^\infty(\Omega)$ and $\psi \in C_c^\infty(\mathbb{R}^d)$, we have

$$\begin{aligned} & \mathbb{E} \int_{\mathbb{R}^d} |\xi(\cdot)| |\psi(X_t^\varepsilon(x)) - \psi(X_t(x))| d\gamma_d(x) \\ & \leq \|\xi\|_\infty \left(\int_{B(R)} + \int_{B(R)^c} \right) \mathbb{E} |\psi(X_t^\varepsilon(x)) - \psi(X_t(x))| d\gamma_d(x) \\ & =: J_1 + J_2. \end{aligned} \quad (3.16)$$

By (3.15),

$$\begin{aligned} J_1 & \leq \|\xi\|_\infty \|\nabla \psi\|_\infty \int_{B(R)} \mathbb{E} |X_t^\varepsilon(x) - X_t(x)| d\gamma_d(x) \\ & \leq \|\xi\|_\infty \|\nabla \psi\|_\infty \left[\sup_{|x| \leq R} \mathbb{E} \left(\sup_{0 \leq t \leq T} |X_t^\varepsilon(x) - X_t(x)|^2 \right) \right]^{1/2} \rightarrow 0, \end{aligned} \quad (3.17)$$

as ε tends to 0. It is obvious that

$$J_2 \leq 2 \|\xi\|_\infty \|\psi\|_\infty \gamma_d(B(R)^c). \quad (3.18)$$

Combining (3.16), (3.17) and (3.18), we obtain

$$\limsup_{\varepsilon \rightarrow 0} \mathbb{E} \int_{\mathbb{R}^d} |\xi| |\psi(X_t^\varepsilon(x)) - \psi(X_t(x))| d\gamma_d(x) \leq 2 \|\xi\|_\infty \|\psi\|_\infty \gamma_d(B(R)^c) \rightarrow 0$$

as $R \uparrow \infty$. Therefore

$$\lim_{\varepsilon \rightarrow 0} \mathbb{E} \int_{\mathbb{R}^d} \xi \psi(X_t^\varepsilon(x)) d\gamma_d(x) = \mathbb{E} \int_{\mathbb{R}^d} \xi \psi(X_t(x)) d\gamma_d(x). \quad (3.19)$$

On the other hand, by Theorem 3.3, for each fixed $t \in [0, T]$, up to a subsequence, K_t^ε converges weakly in $L^1(\Omega \times \mathbb{R}^d)$ to some \hat{K}_t , hence

$$\begin{aligned} \mathbb{E} \int_{\mathbb{R}^d} \xi \psi(X_t^\varepsilon(x)) d\gamma_d(x) & = \mathbb{E} \int_{\mathbb{R}^d} \xi \psi(y) K_t^\varepsilon(y) d\gamma_d(y) \\ & \rightarrow \mathbb{E} \int_{\mathbb{R}^d} \xi \psi(y) \hat{K}_t(y) d\gamma_d(y). \end{aligned} \quad (3.20)$$

This together with (3.19) leads to

$$\mathbb{E} \int_{\mathbb{R}^d} \xi \psi(X_t(x)) d\gamma_d(x) = \mathbb{E} \int_{\mathbb{R}^d} \xi \psi(y) \hat{K}_t(y) d\gamma_d(y).$$

By the arbitrariness of $\xi \in L^\infty(\Omega)$, there exists a full measure subset Ω_ψ of Ω such that

$$\int_{\mathbb{R}^d} \psi(X_t(x)) d\gamma_d(x) = \int_{\mathbb{R}^d} \psi(y) \hat{K}_t(y) d\gamma_d(y), \quad \text{for any } \omega \in \Omega_\psi.$$

Now by the separability of $C_c^\infty(\mathbb{R}^d)$, there exists a full subset Ω_t such that the above equality holds for any $\psi \in C_c^\infty(\mathbb{R}^d)$. Hence $(X_t)_\# \gamma_d = \hat{K}_t \gamma_d$. \square

Remark 3.6. *The $K_t(w, x)$ appearing in (3.13) is defined almost everywhere. It is easy to see that $K_t(w, x)$ is a measurable modification of $\{\hat{K}_t(w, x); t \in [0, T]\}$.*

Remark 3.7. Beyond the Lipschitz condition, several sufficient conditions guaranteeing pathwise uniqueness for SDE (1.1) can be found in the literature. For example in [13], the authors give the condition

$$\sum_{j=1}^m |A_j(x) - A_j(y)|^2 \leq C|x - y|^2 r(|x - y|^2), \quad |A_0(x) - A_0(y)| \leq C|x - y|r(|x - y|^2),$$

for $|x - y| \leq c_0$ small enough, where $r:]0, c_0] \rightarrow]0, +\infty[$ is C^1 satisfying

- (i) $\lim_{s \rightarrow 0} r(s) = +\infty$,
- (ii) $\lim_{s \rightarrow 0} \frac{sr'(s)}{r(s)} = 0$, and
- (iii) $\int_0^{c_0} \frac{ds}{sr(s)} = +\infty$.

Corollary 3.8. Suppose that the vector fields A_0, A_1, \dots, A_m are globally Lipschitz continuous and there exists a constant $C > 0$, such that

$$\sum_{j=1}^m \langle x, A_j(x) \rangle^2 \leq C(1 + |x|^2) \quad \text{for all } x \in \mathbb{R}^d. \quad (3.21)$$

Then $(X_t)_\# \text{Leb}_d \ll \text{Leb}_d$ for any $t \in [0, T]$.

Proof. It is obvious that hypotheses (A1), (A2) and (A4) are satisfied, and that for some constant $C > 0$,

$$|\delta(A_0)|(x) \leq C(1 + |x|^2).$$

Hence there exists $\lambda_0 > 0$ such that $\int_{\mathbb{R}^d} \exp(\lambda_0 |\delta(A_0)|) d\gamma_d < +\infty$. Finally we have

$$\sum_{j=1}^m |\delta(A_j)|^2(x) \leq 2 \sum_{j=1}^m \langle x, A_j(x) \rangle^2 + 2 \sum_{j=1}^m \text{Lip}(A_j)^2.$$

Therefore, under condition (3.21), there exists $\lambda_0 > 0$ such that

$$\int_{\mathbb{R}^d} \exp \left(\lambda_0 \sum_{j=1}^m |\delta(A_j)|^2 \right) d\gamma_d < +\infty.$$

Hence, hypothesis (A3) is satisfied as well. By Theorem 3.4, we have $(X_t)_\# \gamma_d = \hat{K}_t \gamma_d$. Let A be a Borel subset of \mathbb{R}^d such that $\text{Leb}_d(A) = 0$, then $\gamma_d(A) = 0$; therefore $\int_{\mathbb{R}^d} \mathbf{1}_{\{X_t(x) \in A\}} d\gamma_d(x) = 0$. It follows that $\mathbf{1}_{\{X_t(x) \in A\}} = 0$ for Leb_d almost every x , which implies $\text{Leb}_d(X_t \in A) = 0$; this means that $(X_t)_\# \text{Leb}_d$ is absolutely continuous with respect to Leb_d . \square

In the next section, we shall prove that under the conditions of Corollary 3.8, the density of $(X_t)_\# \text{Leb}_d$ with respect to Leb_d is strictly positive, in other words, Leb_d is quasi-invariant under X_t .

Corollary 3.9. Assume that conditions (A1)–(A4) hold. Let $\sigma = (A_j^i)$ and suppose that for some $C > 0$,

$$\sigma(x)\sigma(x)^* \geq C \text{Id}, \quad \text{for all } x \in \mathbb{R}^d.$$

Then $(X_t)_\# \gamma_d$ is absolutely continuous with respect to γ_d .

Proof. The conditions (A1)–(A4) are stronger than those in Theorem 1.1 of [34] given by X. Zhang, so the pathwise uniqueness holds. Hence Theorem 3.4 applies to this case. \square

4 Quasi-invariance under stochastic flow

In the sequel, by quasi-invariance we mean that the Radon-Nikodym derivative of the corresponding push-forward measure is strictly positive. First we prove that in the situation of Corollary 3.8, the Lebesgue measure is in fact quasi-invariant under the stochastic flow of homeomorphisms. To this end, we need some preparations. In what follows, $T_0 > 0$ is chosen small enough such that (3.5) holds.

Proposition 4.1. *Let $q \geq 2$. Then*

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^d} \mathbb{E} \left(\left| \sup_{0 \leq t \leq T_0} \sum_{j=1}^m \int_0^t [\delta(A_j^\varepsilon)(X_s^\varepsilon) - \delta(A_j)(X_s)] dw_s^j \right|^q \right) d\gamma_d = 0. \quad (4.1)$$

Proof. By Burkholder's inequality,

$$\begin{aligned} & \mathbb{E} \left(\sup_{0 \leq t \leq T_0} \left| \sum_{j=1}^m \int_0^t [\delta(A_j^\varepsilon)(X_s^\varepsilon) - \delta(A_j)(X_s)] dw_s^j \right|^q \right) \\ & \leq C \mathbb{E} \left[\left(\int_0^{T_0} \sum_{j=1}^m |\delta(A_j^\varepsilon)(X_s^\varepsilon) - \delta(A_j)(X_s)|^2 ds \right)^{q/2} \right] \\ & \leq C T_0^{q/2-1} \sum_{j=1}^m \int_0^{T_0} \mathbb{E}(|\delta(A_j^\varepsilon)(X_s^\varepsilon) - \delta(A_j)(X_s)|^q) ds. \end{aligned}$$

Again by the inequality $(a+b)^q \leq C_q (a^q + b^q)$, there exists a constant $C_{q,T_0} > 0$ such that the above quantity is dominated by

$$C_{q,T_0} \sum_{j=1}^m \left[\int_0^{T_0} \mathbb{E}(|\delta(A_j^\varepsilon)(X_s^\varepsilon) - \delta(A_j)(X_s)|^q) ds + \int_0^{T_0} \mathbb{E}(|\delta(A_j)(X_s^\varepsilon) - \delta(A_j)(X_s)|^q) ds \right]. \quad (4.2)$$

Let I_1^ε and I_2^ε be the two terms in the squared bracket of (4.2). Note that

$$\begin{aligned} & \int_{\mathbb{R}^d} \mathbb{E}(|\delta(A_j^\varepsilon)(X_s^\varepsilon) - \delta(A_j)(X_s)|^q) d\gamma_d \\ & = \mathbb{E} \int_{\mathbb{R}^d} |\delta(A_j^\varepsilon) - \delta(A_j)|^q K_s^\varepsilon d\gamma_d \\ & \leq \|\delta(A_j^\varepsilon) - \delta(A_j)\|_{L^{2q}(\gamma_d)}^q \|K_s^\varepsilon\|_{L^2(\mathbb{P} \times \gamma_d)}. \end{aligned} \quad (4.3)$$

According to (3.5), for $s \leq T_0$, we have $\|K_s^\varepsilon\|_{L^2(\mathbb{P} \times \gamma_d)} \leq \Lambda_{T_0}$. Remark that

$$\delta(A_j^\varepsilon) = \delta(\varphi_\varepsilon P_\varepsilon A_j) = \varphi_\varepsilon e^\varepsilon P_\varepsilon \delta(A_j) - \langle \nabla \varphi_\varepsilon, P_\varepsilon A_j \rangle,$$

which converges to $\delta(A_j)$ in $L^{2q}(\gamma_d)$. By (4.3),

$$\begin{aligned} \int_{\mathbb{R}^d} I_1^\varepsilon d\gamma_d & = \int_0^{T_0} \left[\int_{\mathbb{R}^d} \mathbb{E}(|\delta(A_j^\varepsilon)(X_s^\varepsilon) - \delta(A_j)(X_s)|^q) d\gamma_d \right] ds \\ & \leq T_0 \Lambda_{T_0} \|\delta(A_j^\varepsilon) - \delta(A_j)\|_{L^{2q}(\gamma_d)}^q \end{aligned}$$

which tends to 0 as $\varepsilon \rightarrow 0$.

For the estimate of I_2^ε , we remark that $\int_{\mathbb{R}^d} |\delta(A_j)|^{2q} d\gamma_d < +\infty$. Let $\eta > 0$ be given. There exists $\psi \in C_c(\mathbb{R}^d)$ such that

$$\int_{\mathbb{R}^d} |\delta(A_j) - \psi|^{2q} d\gamma_d \leq \eta^2.$$

We have, for some constant $C_q > 0$,

$$\begin{aligned} & \int_{\mathbb{R}^d} \mathbb{E}(|\delta(A_j)(X_s^\varepsilon) - \delta(A_j)(X_s)|^q) d\gamma_d \\ & \leq C_q \left[\int_{\mathbb{R}^d} \mathbb{E}(|\delta(A_j)(X_s^\varepsilon) - \psi(X_s)|^q) d\gamma_d + \int_{\mathbb{R}^d} \mathbb{E}(|\psi(X_s^\varepsilon) - \psi(X_s)|^q) d\gamma_d \right. \\ & \quad \left. + \int_{\mathbb{R}^d} \mathbb{E}(|\psi(X_s) - \delta(A_j)(X_s)|^q) d\gamma_d \right]. \end{aligned} \quad (4.4)$$

Again by (3.6), we find

$$\begin{aligned} \mathbb{E} \left[\int_{\mathbb{R}^d} |\delta(A_j)(X_s^\varepsilon) - \psi(X_s)|^q d\gamma_d \right] &= \mathbb{E} \left[\int_{\mathbb{R}^d} |\delta(A_j) - \psi|^q K_s^\varepsilon d\gamma_d \right] \\ &\leq \|\delta(A_j) - \psi\|_{L^{2q}(\gamma_d)}^q \Lambda_{T_0} \leq \Lambda_{T_0} \eta. \end{aligned}$$

In the same way,

$$\mathbb{E} \left[\int_{\mathbb{R}^d} |\delta(A_j)(X_s) - \psi(X_s)|^q d\gamma_d \right] \leq \Lambda_{T_0} \eta.$$

To estimate the second term on the right hand side of (4.4), we use Theorem 3.5: from (3.14), we see that up to a subsequence, $X_s^\varepsilon(w, x)$ converges to $X_s(w, x)$, for each $s \leq T_0$ and almost all $(w, x) \in \Omega \times \mathbb{R}^d$. By Lebesgue's dominated convergence theorem,

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^d} \mathbb{E}(|\psi(X_s^\varepsilon) - \psi(X_s)|^q) d\gamma_d = 0.$$

In conclusion, $\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^d} I_2^\varepsilon d\gamma_d = 0$. According to (4.2), the proof of (4.1) is complete. \square

Proposition 4.2. *Let Φ be defined by*

$$\Phi = \delta(A_0) + \frac{1}{2} \sum_{j=1}^m |A_j|^2 + \frac{1}{2} \sum_{j=1}^m \langle \nabla A_j, (\nabla A_j)^* \rangle, \quad (4.5)$$

and analogously Φ_ε where A_j is replaced by A_j^ε . Then

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^d} \int_0^{T_0} \mathbb{E}(|\Phi_\varepsilon(X_s^\varepsilon) - \Phi(X_s)|^q) ds d\gamma_d = 0. \quad (4.6)$$

Proof. Along the lines of the proof of Proposition 4.1, it is sufficient to remark that

$$\lim_{\varepsilon \rightarrow 0} \|\Phi_\varepsilon - \Phi\|_{L^{2q}(\gamma_d)} = 0. \quad (4.7)$$

To see this, let us check convergence for the last term in the definition of Φ_ε . We have

$$\begin{aligned} & |\langle \nabla A_j^\varepsilon, (\nabla A_j^\varepsilon)^* \rangle - \langle \nabla A_j, (\nabla A_j)^* \rangle| \\ & \leq \|\nabla A_j^\varepsilon - \nabla A_j\| \|\nabla A_j^\varepsilon\| + \|\nabla A_j\| \|\nabla A_j^\varepsilon - \nabla A_j\|. \end{aligned}$$

Note that $A_j^\varepsilon = \varphi_\varepsilon P_\varepsilon A_j$. Thus

$$\nabla A_j^\varepsilon = \nabla \varphi_\varepsilon \otimes P_\varepsilon A_j + e^{-\varepsilon} \varphi_\varepsilon P_\varepsilon (\nabla A_j),$$

which converges to ∇A_j in $L^{2q}(\gamma_d)$ as $\varepsilon \rightarrow 0$. \square

Now we can prove

Proposition 4.3. *Under the conditions of Corollary 3.8, the Lebesgue measure Leb_d is quasi-invariant under the stochastic flow.*

Proof. Let k_t be the density of $(X_t)_\# \text{Leb}_d$ with respect to Leb_d . We shall prove that k_t is strictly positive. Set

$$\tilde{K}_t^\varepsilon(x) = \exp \left(- \sum_{j=1}^m \int_0^t \delta(A_j^\varepsilon)(X_s^\varepsilon(x)) dw_s^j - \int_0^t \Phi_\varepsilon(X_s^\varepsilon(x)) ds \right), \quad (4.8)$$

where Φ_ε is defined in Proposition 4.2. By (2.3) we have

$$\int_{\mathbb{R}^d} \psi(X_t^\varepsilon) \tilde{K}_t^\varepsilon d\gamma_d = \int_{\mathbb{R}^d} \psi d\gamma_d, \quad \psi \in C_c^1(\mathbb{R}^d). \quad (4.9)$$

Applying Propositions 4.1 and 4.2, up to a subsequence, for each $t \leq T_0$ and almost every (w, x) , the term $\tilde{K}_t^\varepsilon(w, x)$ defined in (4.8) converges to

$$\tilde{K}_t(x) = \exp \left(- \sum_{j=1}^m \int_0^t \delta(A_j)(X_s(x)) dw_s^j - \int_0^t \Phi(X_s(x)) ds \right). \quad (4.10)$$

By Corollary 2.3 and Lemma 3.2, we may assume that T_0 is small enough so that for any $t \leq T_0$, the family $\{\tilde{K}_t^\varepsilon : \varepsilon \leq 1\}$ is also bounded in $L^2(\mathbb{P} \times \gamma_d)$. Therefore, by the uniform integrability, letting $\varepsilon \rightarrow 0$ in (4.9), we get \mathbb{P} -almost surely,

$$\int_{\mathbb{R}^d} \psi(X_t) \tilde{K}_t d\gamma_d = \int_{\mathbb{R}^d} \psi d\gamma_d, \quad \psi \in C_c^1(\mathbb{R}^d). \quad (4.11)$$

Now taking a Borel version of $x \rightarrow \tilde{K}_t(w, x)$. Under the assumptions, the solution X_t is a stochastic flow of homeomorphisms, hence the inverse flow X_t^{-1} exists. Consequently, if $t \leq T_0$, we deduce from (4.11) that the density $K_t(w, x)$ of $(X_t)_\# \gamma_d$ with respect to γ_d admits the expression $K_t(w, x) = [\tilde{K}_t(w, X_t^{-1}(w, x))]^{-1}$ which is strictly positive. For X_{t+T_0} with $t \leq T_0$, we use the flow property: $X_{t+T_0}(w, x) = X_t(\theta_{T_0}w, X_{T_0}(w, x))$. Thus, for any $\psi \in C_c^1(\mathbb{R}^d)$,

$$\begin{aligned} \int_{\mathbb{R}^d} \psi(X_{t+T_0}) d\gamma_d &= \int_{\mathbb{R}^d} \psi(X_t(X_{T_0})) d\gamma_d \\ &= \int_{\mathbb{R}^d} \psi(X_t) K_{T_0} d\gamma_d = \int_{\mathbb{R}^d} \psi K_{T_0}(X_t^{-1}) K_t d\gamma_d. \end{aligned}$$

That is to say, the density $K_{t+T_0} = K_{T_0}(X_t^{-1}) K_t$ is strictly positive. Continuing in this way, we obtain that K_t is strictly positive for any $t \geq 0$.

Now if $\rho(x)$ denotes the density of γ_d with respect to Leb_d , then

$$k_t(w, x) = \rho(X_t^{-1}(w, x))^{-1} K_t(w, x) \rho(x) > 0$$

which concludes the proof. \square

In what follows, we will give examples for which existence of the inverse flow is not known.

Theorem 4.4. *Let A_1, \dots, A_m be bounded C^1 vector fields on \mathbb{R}^d such that their derivatives are of linear growth; furthermore let A_0 be continuous of linear growth such that $\delta(A_0)$ exists. Define*

$$\hat{A}_0 = A_0 - \sum_{j=1}^m \mathcal{L}_{A_j} A_j. \quad (4.12)$$

Suppose that $\delta(\hat{A}_0)$ exists and that

$$\int_{\mathbb{R}^d} \exp(\lambda_0 (|\delta(A_0)| + |\delta(\hat{A}_0)|)) d\gamma_d < +\infty, \quad \text{for some } \lambda_0 > 0. \quad (4.13)$$

If pathwise uniqueness holds both for SDE (1.1) and for

$$dY_t = \sum_{j=1}^m A_j(Y_t) dw_t^j - \hat{A}_0(Y_t) dt, \quad (4.14)$$

then the solution X_t to SDE (1.1) leaves the Gaussian measure γ_d quasi-invariant.

Proof. Obviously the conditions in Theorem 3.4 are satisfied; hence $(X_t)_{\#}\gamma_d = K_t \gamma_d$. Let $t > 0$ be given, we consider the dual SDE to (1.1):

$$dY_s^t = \sum_{j=1}^m A_j(Y_s^t) dw_s^{t,j} - \hat{A}_0(Y_s^t) ds$$

for which pathwise uniqueness holds; here $w_s^t = w_{t-s} - w_t$ with $s \in [0, t]$. Let A_j^ε , $j = 0, 1, \dots, m$, be the vector fields defined as above. Consider

$$dY_s^{t,\varepsilon} = \sum_{j=1}^m A_j^\varepsilon(Y_s^{t,\varepsilon}) dw_s^{t,j} - \hat{A}_0^\varepsilon(Y_s^{t,\varepsilon}) ds,$$

where $\hat{A}_0^\varepsilon = A_0^\varepsilon - \sum_{j=1}^m \mathcal{L}_{A_j^\varepsilon} A_j^\varepsilon$. Then it is known that $(X_t^\varepsilon)^{-1} = Y_t^{t,\varepsilon}$. It is easy to check that for some constant $C > 0$ independent of ε ,

$$|\hat{A}_0^\varepsilon(x)| \leq C(1 + |x|). \quad (4.15)$$

Moreover,

$$\mathcal{L}_{A_j^\varepsilon} A_j^\varepsilon = \sum_{k=1}^d (A_j^\varepsilon)^k \left[\frac{\partial \varphi_\varepsilon}{\partial x_k} P_\varepsilon A_j + \varphi_\varepsilon e^{-\varepsilon} P_\varepsilon \left(\frac{\partial A_j}{\partial x_k} \right) \right]$$

which converges locally uniformly to $\mathcal{L}_{A_j} A_j$. Therefore \hat{A}_0^ε converges uniformly over any compact subset to \hat{A}_0 . By Theorem 3.5,

$$\lim_{\varepsilon \rightarrow 0} \sup_{|x| \leq R} \mathbb{E} \left(\sup_{0 \leq s \leq t} |Y_s^{t,\varepsilon} - Y_s^t|^2 \right) = 0.$$

It follows that, along a sequence, $Y_t^{t,\varepsilon}$ converges to Y_t^t for almost every (w, x) . Now let $\psi_1, \psi_2 \in C_b(\mathbb{R}^d)$, we have for $t \leq T_0$,

$$\int_{\mathbb{R}^d} \psi_1 \cdot \psi_2(X_t^\varepsilon) \tilde{K}_t^\varepsilon d\gamma_d = \int_{\mathbb{R}^d} \psi_1(Y_t^{t,\varepsilon}) \cdot \psi_2 d\gamma_d.$$

Letting $\varepsilon \rightarrow 0$ leads to

$$\int_{\mathbb{R}^d} \psi_1 \cdot \psi_2(X_t) \tilde{K}_t d\gamma_d = \int_{\mathbb{R}^d} \psi_1(Y_t^t) \cdot \psi_2 d\gamma_d. \quad (4.16)$$

Taking ψ_1 and ψ_2 positive in (4.16) and using a monotone class argument, we see that equation (4.16) holds for any positive Borel functions ψ_1 and ψ_2 . Hence taking a Borel version of \tilde{K}_t and setting $\psi_1 = 1/\tilde{K}_t$ in (4.16), we get

$$\int_{\mathbb{R}^d} \psi_2(X_t) d\gamma_d = \int_{\mathbb{R}^d} [\tilde{K}_t(Y_t^t)]^{-1} \psi_2 d\gamma_d. \quad (4.17)$$

It follows that $K_t = [\tilde{K}_t(Y_t^t)]^{-1} > 0$ for $t \leq T_0$. For X_{t+T_0} with $t \leq T_0$, we shall use repeatedly (4.16). By the flow property, $X_{t+T_0}(w, x) = X_t(\theta_{T_0}w, X_{T_0}(w, x))$ where $(\theta_{T_0}w)_t = w_{t+T_0} - w_{T_0}$. Letting $t = T_0$ and replacing ψ_2 by $\psi_2(X_t)$ we get

$$\int_{\mathbb{R}^d} \psi_1 \cdot \psi_2(X_{t+T_0}) \tilde{K}_{T_0} d\gamma_d = \int_{\mathbb{R}^d} \psi_1(Y_{T_0}^{T_0}) \psi_2(X_t) d\gamma_d.$$

Taking $\psi_1 = 1/\tilde{K}_{T_0}$ in the above equality, we get

$$\begin{aligned} \int_{\mathbb{R}^d} \psi_2(X_{t+T_0}) d\gamma_d &= \int_{\mathbb{R}^d} [\tilde{K}_{T_0}(Y_{T_0}^{T_0})]^{-1} \psi_2(X_t) d\gamma_d \\ &= \int_{\mathbb{R}^d} [\tilde{K}_{T_0}(Y_{T_0}^{T_0})]^{-1} \psi_2(X_t) \tilde{K}_t^{-1} \tilde{K}_t d\gamma_d \\ &= \int_{\mathbb{R}^d} [\tilde{K}_{T_0}(Y_{T_0}^{T_0}(Y_t^t))]^{-1} [\tilde{K}_t(Y_t^t)]^{-1} \psi_2 d\gamma_d, \end{aligned}$$

where in the last equality we have used (4.16) with $\psi_1 = [\tilde{K}_{T_0}(Y_{T_0}^{T_0})]^{-1} \tilde{K}_t^{-1}$. It follows that the density K_{t+T_0} of $(X_{t+T_0})_{\#} \gamma_d$ with respect to γ_d is strictly positive, and so on. \square

Corollary 4.5. *Let A_1, \dots, A_m be bounded C^2 vector fields such that their derivatives up to order 2 grow at most linearly, and let A_0 be a continuous vector field of linear growth. Suppose that*

$$|A_0(x) - A_0(y)| \leq C_R |x - y| \log_k \frac{1}{|x - y|} \quad \text{for } |x| \leq R, |y| \leq R, |x - y| \leq c_0 \text{ small enough,} \quad (4.18)$$

where $\log_k s = (\log s)(\log \log s) \dots (\log \dots \log s)$. Suppose further that

$$\text{div}(A_0) = \sum_{j=1}^d \frac{\partial A_0^j}{\partial x_j}$$

exists and is bounded. Then the stochastic flow X_t defined by SDE (1.1) leaves the Lebesgue measure quasi-invariant.

Proof. It is obvious that \hat{A}_0 defined in (4.12) satisfies condition (4.18); therefore by [13], pathwise uniqueness holds for SDE (1.1) and (4.14). Note that $\delta(A_0) = \langle x, A_0 \rangle - \text{div}(A_0)$. Then condition (4.13) is satisfied; thus Theorem 4.4 yields the result. \square

5 The case A_0 in Sobolev spaces

From now on, A_0 is not supposed to be continuous, but in some Sobolev space, that is, we replace the condition (A1) in **(H)** by

(A1') For $i = 1, \dots, m$, $A_i \in \cap_{q \geq 1} \mathbb{D}_1^q(\gamma_d)$, $A_0 \in \mathbb{D}_1^q(\gamma_d)$ for some $q > 1$.

First we establish the following *a priori* estimate on perturbations, using the method developed in [36]. Let $\{A_0, A_1, \dots, A_m\}$ be a family of measurable vector fields on \mathbb{R}^d . We shall give a precise definition of solution to the following SDE

$$dX_t = \sum_{i=1}^m A_i(X_t) dw_t^i + A_0(X_t) dt, \quad X_0 = x. \quad (5.1)$$

Definition 5.1. We say that a measurable map $X : \Omega \times \mathbb{R}^d \rightarrow C([0, T], \mathbb{R}^d)$ is a solution to Itô SDE (5.1) if

- (i) for each $t \in [0, T]$ and almost all $x \in \mathbb{R}^d$, $w \rightarrow X_t(w, x)$ is measurable with respect to \mathcal{F}_t , i.e. the natural filtration generated by the Brownian motion $\{w_s; s \leq t\}$;
- (ii) for each $t \in [0, T]$, there exists $K_t \in L^1(\mathbb{P} \times \mathbb{R}^d)$ such that $(X_t(w, \cdot))_{\#} \gamma_d$ admits K_t as the density with respect to γ_d ;
- (iii) almost surely

$$\sum_{i=1}^m \int_0^T |A_i(X_s(w, x))|^2 ds + \int_0^T |A_0(X_s(w, x))| ds < +\infty;$$

- (iv) for almost all $x \in \mathbb{R}^d$,

$$X_t(w, x) = x + \sum_{i=1}^m \int_0^t A_i(X_s(w, x)) dw_s^i + \int_0^t A_0(X_s(w, x)) ds;$$

- (v) the flow property holds

$$X_{t+s}(w, x) = X_t(\theta_s w, X_s(w, x)).$$

Now consider another family of measurable vector fields $\{\hat{A}_0, \hat{A}_1, \dots, \hat{A}_m\}$ on \mathbb{R}^d , and denote by \hat{X}_t the solution to the SDE

$$d\hat{X}_t = \sum_{i=1}^m \hat{A}_i(\hat{X}_t) dw_t^i + \hat{A}_0(\hat{X}_t) dt, \quad \hat{X}_0 = x. \quad (5.2)$$

Let \hat{K}_t be the density of $(\hat{X}_t)_{\#} \gamma_d$ and define

$$\Lambda_{p,T} = \sup_{0 \leq t \leq T} \left(\|K_t\|_{L^p(\mathbb{P} \times \gamma_d)} \vee \|\hat{K}_t\|_{L^p(\mathbb{P} \times \gamma_d)} \right). \quad (5.3)$$

Theorem 5.2. Let $q > 1$. Suppose that A_1, \dots, A_m as well as $\hat{A}_1, \dots, \hat{A}_m$ are in $\mathbb{D}_1^{2q}(\gamma_d)$ and $A_0, \hat{A}_0 \in \mathbb{D}_1^q(\gamma_d)$. Then for any $T > 0$ and $R > 0$, there exist constants $C_{d,q,R} > 0$ and $C_T > 0$ such that for any $\sigma > 0$,

$$\begin{aligned} & \mathbb{E} \left[\int_{G_R} \log \left(\frac{\sup_{0 \leq t \leq T} |X_t - \hat{X}_t|^2}{\sigma^2} + 1 \right) d\gamma_d \right] \\ & \leq C_T \Lambda_{p,T} \left\{ C_{d,q,R} \left[\|\nabla A_0\|_{L^q} + \left(\sum_{i=1}^m \|\nabla A_i\|_{L^{2q}}^2 \right)^{\frac{1}{2}} + \sum_{i=1}^m \|\nabla A_i\|_{L^{2q}}^2 \right] \right. \\ & \quad \left. + \frac{1}{\sigma^2} \sum_{i=1}^m \|A_i - \hat{A}_i\|_{L^{2q}}^2 + \frac{1}{\sigma} \left[\|A_0 - \hat{A}_0\|_{L^q} + \left(\sum_{i=1}^m \|A_i - \hat{A}_i\|_{L^{2q}}^2 \right)^{\frac{1}{2}} \right] \right\}, \end{aligned}$$

where p is the conjugate number of q : $1/p + 1/q = 1$ and

$$G_R(w) = \left\{ x \in \mathbb{R}^d : \sup_{0 \leq t \leq T} |X_t(w, x)| \vee |\hat{X}_t(w, x)| \leq R \right\}. \quad (5.4)$$

Proof. Denote by $\xi_t = X_t - \hat{X}_t$, then $\xi_0 = 0$. By Itô formula,

$$\begin{aligned} d|\xi_t|^2 &= 2 \sum_{i=1}^m \langle \xi_t, A_i(X_t) - \hat{A}_i(\hat{X}_t) \rangle dw_t^i + 2 \langle \xi_t, A_0(X_t) - \hat{A}_0(\hat{X}_t) \rangle dt \\ &\quad + \sum_{i=1}^m |A_i(X_t) - \hat{A}_i(\hat{X}_t)|^2 dt. \end{aligned} \quad (5.5)$$

For $\sigma > 0$, $\log\left(\frac{|\xi_t|^2}{\sigma^2} + 1\right) = \log(|\xi_t|^2 + \sigma^2) - \log\sigma^2$. Again by the Itô formula,

$$d\log(|\xi_t|^2 + \sigma^2) = \frac{d|\xi_t|^2}{|\xi_t|^2 + \sigma^2} - \frac{1}{2} \cdot \frac{4 \sum_{i=1}^m \langle \xi_t, A_i(X_t) - \hat{A}_i(\hat{X}_t) \rangle^2}{(|\xi_t|^2 + \sigma^2)^2} dt,$$

using (5.5), we obtain

$$\begin{aligned} d\log(|\xi_t|^2 + \sigma^2) &= 2 \sum_{i=1}^m \frac{\langle \xi_t, A_i(X_t) - \hat{A}_i(\hat{X}_t) \rangle}{|\xi_t|^2 + \sigma^2} dw_t^i + 2 \frac{\langle \xi_t, A_0(X_t) - \hat{A}_0(\hat{X}_t) \rangle}{|\xi_t|^2 + \sigma^2} dt \\ &\quad + \sum_{i=1}^m \frac{|A_i(X_t) - \hat{A}_i(\hat{X}_t)|^2}{|\xi_t|^2 + \sigma^2} dt - 2 \sum_{i=1}^m \frac{\langle \xi_t, A_i(X_t) - \hat{A}_i(\hat{X}_t) \rangle^2}{(|\xi_t|^2 + \sigma^2)^2} dt \\ &=: dI_1(t) + dI_2(t) + dI_3(t) + dI_4(t). \end{aligned} \quad (5.6)$$

Let $\tau_R(x) = \inf\{t \geq 0 : |X_t(x)| \vee |\hat{X}_t(x)| > R\}$. Remark that almost surely, $G_R \subset \{x : \tau_R(x) > T\}$ and for any $t \geq 0$, $\{\tau_R > t\} \subset B(R)$. Therefore

$$\mathbb{E} \left[\int_{G_R} \sup_{0 \leq t \leq T} |I_1(t)| d\gamma_d \right] \leq \mathbb{E} \left[\int_{B(R)} \sup_{0 \leq t \leq T \wedge \tau_R} |I_1(t)| d\gamma_d \right].$$

By Burkholder's inequality,

$$\mathbb{E} \left(\sup_{0 \leq t \leq T \wedge \tau_R} |I_1(t)|^2 \right) \leq 4 \mathbb{E} \left(\int_0^{T \wedge \tau_R} \sum_{i=1}^m \frac{\langle \xi_t, A_i(X_t) - \hat{A}_i(\hat{X}_t) \rangle^2}{(|\xi_t|^2 + \sigma^2)^2} dt \right),$$

which is obviously less than

$$4 \mathbb{E} \left(\int_0^{T \wedge \tau_R} \sum_{i=1}^m \frac{|A_i(X_t) - \hat{A}_i(\hat{X}_t)|^2}{|\xi_t|^2 + \sigma^2} dt \right).$$

Hence

$$\mathbb{E} \left[\int_{B(R)} \sup_{0 \leq t \leq T \wedge \tau_R} |I_1(t)| d\gamma_d \right] \leq 4 \left[\int_0^T \left(\mathbb{E} \int_{\{\tau_R > t\}} \sum_{i=1}^m \frac{|A_i(X_t) - \hat{A}_i(\hat{X}_t)|^2}{|\xi_t|^2 + \sigma^2} d\gamma_d \right) dt \right]^{\frac{1}{2}}. \quad (5.7)$$

We have $A_i(X_t) - \hat{A}_i(\hat{X}_t) = A_i(X_t) - A_i(\hat{X}_t) + A_i(\hat{X}_t) - \hat{A}_i(\hat{X}_t)$. Using the density \hat{K}_t , it is clear that

$$\begin{aligned} \mathbb{E} \int_{\{\tau_R > t\}} \frac{|A_i(\hat{X}_t) - \hat{A}_i(\hat{X}_t)|^2}{|\xi_t|^2 + \sigma^2} d\gamma_d &\leq \frac{1}{\sigma^2} \mathbb{E} \int_{\mathbb{R}^d} |A_i(\hat{X}_t) - \hat{A}_i(\hat{X}_t)|^2 d\gamma_d \\ &= \frac{1}{\sigma^2} \mathbb{E} \int_{\mathbb{R}^d} |A_i - \hat{A}_i|^2 \hat{K}_t d\gamma_d. \end{aligned}$$

Thus by Hölder's inequality and according to (5.3), we have

$$\mathbb{E} \int_{\{\tau_R > t\}} \frac{|A_i(\hat{X}_t) - \hat{A}_i(\hat{X}_t)|^2}{|\xi_t|^2 + \sigma^2} d\gamma_d \leq \frac{\Lambda_{p,T}}{\sigma^2} \|A_i - \hat{A}_i\|_{L^{2q}}^2. \quad (5.8)$$

Now we shall use Theorem 6.1 in the Appendix to estimate another term. Note that on the set $\{\tau_R > t\}$, $X_t, \hat{X}_t \in B(R)$, then $|X_t - \hat{X}_t| \leq 2R$. Since $(X_t)_{\#} \gamma_d \ll \gamma_d$ and $(\hat{X}_t)_{\#} \gamma_d \ll \gamma_d$, we can apply (6.2) so that

$$|A_i(X_t) - A_i(\hat{X}_t)| \leq C_d |X_t - \hat{X}_t| (M_{2R} |\nabla A_i|(X_t) + M_{2R} |\nabla A_i|(\hat{X}_t)).$$

Then

$$\mathbb{E} \left[\int_{\{\tau_R > t\}} \frac{|A_i(X_t) - A_i(\hat{X}_t)|^2}{|\xi_t|^2 + \sigma^2} d\gamma_d \right] \leq C_d^2 \mathbb{E} \int_{\{\tau_R > t\}} (M_{2R} |\nabla A_i|(X_t) + M_{2R} |\nabla A_i|(\hat{X}_t))^2 d\gamma_d.$$

Notice again that on $\{\tau_R(x) > t\}$, $X_t(x)$ and $\hat{X}_t(x)$ are in $B(R)$, therefore

$$\begin{aligned} \mathbb{E} \left[\int_{\{\tau_R > t\}} \frac{|A_i(X_t) - A_i(\hat{X}_t)|^2}{|\xi_t|^2 + \sigma^2} d\gamma_d \right] &\leq 2C_d^2 \mathbb{E} \int_{B(R)} (M_{2R} |\nabla A_i|)^2 (K_t + \hat{K}_t) d\gamma_d \\ &\leq 4C_d^2 \Lambda_{p,T} \left(\int_{B(R)} (M_{2R} |\nabla A_i|)^{2q} d\gamma_d \right)^{\frac{1}{q}}. \end{aligned} \quad (5.9)$$

Remark that the maximal function inequality does not hold for the Gaussian measure γ_d on the whole space \mathbb{R}^d . However, on each ball $B(R)$,

$$\gamma_d|_{B(R)} \leq \frac{1}{(2\pi)^{d/2}} \text{Leb}_d|_{B(R)} \leq e^{R^2/2} \gamma_d|_{B(R)}.$$

Thus, according to (6.3),

$$\begin{aligned} \int_{B(R)} (M_{2R} |\nabla A_i|)^{2q} d\gamma_d &\leq \frac{1}{(2\pi)^{d/2}} \int_{B(R)} (M_{2R} |\nabla A_i|)^{2q} dx \leq \frac{C_{d,q}}{(2\pi)^{d/2}} \int_{B(3R)} |\nabla A_i|^{2q} dx \\ &\leq C_{d,q} e^{9R^2/2} \int_{B(3R)} |\nabla A_i|^{2q} d\gamma_d \leq C_{d,q} e^{9R^2/2} \|\nabla A_i\|_{L^{2q}}^{2q}. \end{aligned}$$

Therefore by (5.9), there exists a constant $C_{d,q,R} > 0$ such that

$$\mathbb{E} \left[\int_{\{\tau_R > t\}} \frac{|A_i(X_t) - A_i(\hat{X}_t)|^2}{|\xi_t|^2 + \sigma^2} d\gamma_d \right] \leq C_{d,q,R} \Lambda_{p,T} \|\nabla A_i\|_{L^{2q}}^2.$$

Combining this estimate with (5.7) and (5.8), we get

$$\mathbb{E} \left[\int_{G_R} \sup_{0 \leq t \leq T} |I_1(t)| d\gamma_d \right] \leq CT^{\frac{1}{2}} \Lambda_{p,T}^{\frac{1}{2}} \left(C_{d,q,R} \sum_{i=1}^m \|\nabla A_i\|_{L^{2q}}^2 + \frac{1}{\sigma^2} \sum_{i=1}^m \|A_i - \hat{A}_i\|_{L^{2q}}^2 \right)^{\frac{1}{2}}. \quad (5.10)$$

Now we turn to deal with $I_2(t)$ in (5.6). We have

$$\mathbb{E} \left[\int_{G_R} \sup_{0 \leq t \leq T} |I_2(t)| d\gamma_d \right] \leq 2 \int_0^T \left[\mathbb{E} \int_{G_R} \frac{|A_0(X_t) - \hat{A}_0(\hat{X}_t)|}{(|\xi_t|^2 + \sigma^2)^{\frac{1}{2}}} d\gamma_d \right] dt.$$

Note that for $x \in G_R$, $\hat{X}_t(x) \in B(R)$ for each $t \in [0, T]$, thus

$$\mathbb{E} \left[\int_{G_R} \frac{|A_0(\hat{X}_t) - \hat{A}_0(\hat{X}_t)|}{(|\xi_t|^2 + \sigma^2)^{\frac{1}{2}}} d\gamma_d \right] \leq \frac{1}{\sigma} \mathbb{E} \int_{B(R)} |A_0 - \hat{A}_0| \hat{K}_t d\gamma_d \leq \frac{\Lambda_{p,T}}{\sigma} \|A_0 - \hat{A}_0\|_{L^q}.$$

Again using (6.2),

$$\mathbb{E} \left[\int_{G_R} \frac{|A_0(X_t) - A_0(\hat{X}_t)|}{(|\xi_t|^2 + \sigma^2)^{\frac{1}{2}}} d\gamma_d \right] \leq C_d \mathbb{E} \int_{G_R} (M_{2R} |\nabla A_0|(X_t) + M_{2R} |\nabla A_0|(\hat{X}_t)) d\gamma_d,$$

which is dominated by

$$C_d \mathbb{E} \left[\int_{B(R)} (M_{2R} |\nabla A_0|) \cdot (K_t + \hat{K}_t) d\gamma_d \right] \leq C_{d,q,R} \|\nabla A_0\|_{L^q} \Lambda_{p,T}.$$

Therefore we get the following estimate for I_2 :

$$\mathbb{E} \left[\int_{G_R} \sup_{0 \leq t \leq T} |I_2(t)| d\gamma_d \right] \leq 2T \Lambda_{p,T} \left(C_{d,q,R} \|\nabla A_0\|_{L^q} + \frac{1}{\sigma} \|A_0 - \hat{A}_0\|_{L^q} \right). \quad (5.11)$$

In the same way we have

$$\mathbb{E} \left[\int_{G_R} \sup_{0 \leq t \leq T} |I_3(t)| d\gamma_d \right] \leq CT \Lambda_{p,T} \left(C_{d,q,R} \sum_{i=1}^m \|\nabla A_i\|_{L^{2q}}^2 + \frac{1}{\sigma^2} \sum_{i=1}^m \|A_i - \hat{A}_i\|_{L^{2q}}^2 \right). \quad (5.12)$$

The term $I_4(t)$ is negative and hence we omit it. Combining (5.6) and (5.10)–(5.12), we complete the proof. \square

Now we shall construct a solution to SDE (5.1). To this end, we take $\varepsilon = 1/n$ and we write A_j^n instead of $A_j^{1/n}$ introduced in Section 3. Then by assumption (A2) and Lemma 3.1, there is $C > 0$ independent of n and i , such that

$$|A_i^n(x)| \leq C(1 + |x|). \quad (5.13)$$

Let X_t^n be the solution to Itô SDE (5.1) with the coefficients A_i^n ($i = 0, 1, \dots, m$). Then for any $\alpha \geq 1$ and $T > 0$, there exists $C_{\alpha,T} > 0$ independent of n such that

$$\mathbb{E} \left(\sup_{0 \leq t \leq T} |X_t^n|^\alpha \right) \leq C_{\alpha,T} (1 + |x|^\alpha), \quad \text{for all } x \in \mathbb{R}^d. \quad (5.14)$$

Let K_t^n be the density of $(X_t^n)_{\#}\gamma_d$ with respect to γ_d . Under the hypotheses (A2)–(A4), there is $T_0 > 0$ such that (recall that p is the conjugate number of $q > 1$):

$$\begin{aligned} \Lambda_{p,T_0} := & \left[\int_{\mathbb{R}^d} \exp \left(2pT_0 [|A_0| + e|\delta(A_0)| \right. \right. \\ & \left. \left. + \sum_{j=1}^m (2p|A_j|^2 + |\nabla A_j|^2 + 2(p-1)e^2|\delta(A_j)|^2) \right] \right) d\gamma_d \right]^{\frac{p-1}{p(2p-1)}} < \infty. \end{aligned} \quad (5.15)$$

Similar to (3.6), we have

$$\sup_{t \in [0, T_0]} \sup_{n \geq 1} \|K_t^n\|_{L^p(\gamma_d \times \mathbb{P})} \leq \Lambda_{p,T_0} < +\infty. \quad (5.16)$$

Now we shall prove that the family $\{X_t^n : n \geq 1\}$ is convergent to some stochastic field.

Theorem 5.3. *Let T_0 be given in (5.15). Then under the assumptions (A1') and (A2)–(A4), there exists $X : \Omega \times \mathbb{R}^d \rightarrow C([0, T_0], \mathbb{R}^d)$ such that for any $\alpha \geq 1$,*

$$\lim_{n \rightarrow \infty} \mathbb{E} \left[\int_{\mathbb{R}^d} \left(\sup_{0 \leq t \leq T_0} |X_t^n - X_t|^\alpha \right) d\gamma_d \right] = 0. \quad (5.17)$$

Proof. We shall prove that $\{X^n; n \geq 1\}$ is a Cauchy sequence in $L^\alpha(\Omega \times \mathbb{R}^d; C([0, T_0], \mathbb{R}^d))$. Denote by $\|\cdot\|_{\infty, T_0}$ the uniform norm on $C([0, T_0], \mathbb{R}^d)$, so what we have to prove is

$$\lim_{n, k \rightarrow +\infty} \mathbb{E} \left(\int_{\mathbb{R}^d} \|X^n - X^k\|_{\infty, T_0}^\alpha d\gamma_d \right) = 0. \quad (5.18)$$

First by (5.14), the quantity

$$J_{\alpha, T_0} := \sup_{n \geq 1} \mathbb{E} \left(\int_{\mathbb{R}^d} \|X^n\|_{\infty, T_0}^{2\alpha} d\gamma_d \right) \leq C_{\alpha, T_0} \int_{\mathbb{R}^d} (1 + |x|^{2\alpha}) d\gamma_d \quad (5.19)$$

is obviously finite. Let $R > 0$ and set

$$G_{n, R}(w) = \{x \in \mathbb{R}^d; \|X^n(w, x)\|_{\infty, T_0} \leq R\}.$$

Using (5.19), for any $\alpha \geq 1$ and $R > 0$, we have

$$\sup_{n \geq 1} \mathbb{E}(\gamma_d(G_{n, R}^c)) \leq \frac{J_{\alpha, T_0}}{R^{2\alpha}}.$$

Now by Cauchy-Schwarz inequality

$$\begin{aligned} & \mathbb{E} \left(\int_{G_{n, R}^c \cup G_{k, R}^c} \|X^n - X^k\|_{\infty, T_0}^\alpha d\gamma_d \right) \\ & \leq \left(\mathbb{E}[\gamma_d(G_{n, R}^c \cup G_{k, R}^c)] \right)^{1/2} \cdot \left(\mathbb{E} \int_{\mathbb{R}^d} \|X^n - X^k\|_{\infty, T_0}^{2\alpha} d\gamma_d \right)^{1/2} \\ & \leq \left(\frac{2J_{\alpha, T_0}}{R^{2\alpha}} \right)^{1/2} \cdot (2^{2\alpha} J_{\alpha, T_0})^{1/2}. \end{aligned}$$

Let $\varepsilon > 0$ be given; choose $R > 1$ big enough such that the last quantity in the above inequality is less than ε . Then we have for any $n, k \geq 1$,

$$\mathbb{E} \left(\int_{G_{n, R}^c \cup G_{k, R}^c} \|X^n - X^k\|_{\infty, T_0}^\alpha d\gamma_d \right) \leq \varepsilon. \quad (5.20)$$

Let

$$\sigma_{n, k} = \|A_0^n - A_0^k\|_{L^q} + \left(\sum_{i=1}^m \|A_i^n - A_i^k\|_{L^{2q}}^2 \right)^{1/2},$$

which tends to 0 as $n, k \rightarrow +\infty$ since A_0^n converges to A_0 in $L^q(\gamma_d)$ and A_i^n converges to A_i in $L^{2q}(\gamma_d)$ for $i = 1, \dots, m$. Now applying Theorem 5.2 with A_i and \hat{A}_i being replaced respectively by A_i^n and A_i^k , we get

$$\begin{aligned} I_{n, k} &:= \mathbb{E} \left[\int_{G_{n, R} \cap G_{k, R}} \log \left(\frac{\|X^n - X^k\|_{\infty, T_0}^2}{\sigma_{n, k}^2} + 1 \right) d\gamma_d \right] \\ &\leq C_{T_0} \Lambda_{p, T_0} \left\{ C_{d, q, R} \left[\|\nabla A_0^n\|_{L^q} + \left(\sum_{i=1}^m \|\nabla A_i^n\|_{L^{2q}}^2 \right)^{1/2} + \sum_{i=1}^n \|\nabla A_i^n\|_{L^{2q}}^2 \right] + 2 \right\}. \end{aligned}$$

Recall that $A_i^n = \varphi_{1/n} P_{1/n} A_i$, then $\nabla A_i^n = \nabla \varphi_{1/n} \otimes P_{1/n} A_i + \varphi_{1/n} e^{-1/n} P_{1/n} \nabla A_i$, therefore

$$|\nabla A_i^n| \leq P_{1/n}(|A_i| + |\nabla A_i|).$$

We get the following uniform estimates

$$\|\nabla A_0^n\|_{L^q} \leq \|A_0\|_{\mathbb{D}_1^q}, \quad \|\nabla A_i^n\|_{L^{2q}} \leq \|A_i\|_{\mathbb{D}_1^{2q}}.$$

So the quantity $I_{n,k}$ is uniformly bounded with respect to n, k . Let $\hat{\Pi}$ be the measure on $\Omega \times \mathbb{R}^d$ defined by

$$\int_{\Omega \times \mathbb{R}^d} \psi(w, x) d\hat{\Pi}(w, x) = \mathbb{E} \left[\int_{G_{n,R} \cap G_{k,R}} \psi(w, x) d\gamma_d(x) \right].$$

We have $\hat{\Pi}(\Omega \times \mathbb{R}^d) \leq 1$. Let $\eta > 0$, consider

$$\Sigma_{n,k} = \{(w, x); \|X^n(w, x) - X^k(w, x)\|_{\infty, T_0} \geq \eta\},$$

which is equal to

$$\left\{ (w, x); \log \left(\frac{\|X^n - X^k\|_{\infty, T_0}^2}{\sigma_{n,k}^2} + 1 \right) \geq \log \left(\frac{\eta^2}{\sigma_{n,k}^2} + 1 \right) \right\}.$$

It follows that as $n, k \rightarrow +\infty$,

$$\hat{\Pi}(\Sigma_{n,k}) \leq \frac{I_{n,k}}{\log \left(\frac{\eta^2}{\sigma_{n,k}^2} + 1 \right)} \rightarrow 0, \quad (5.21)$$

since $\sigma_{n,k} \rightarrow 0$ and the family $\{I_{n,k}; n, k \geq 1\}$ is bounded. Now

$$\begin{aligned} \mathbb{E} \left(\int_{G_{n,R} \cap G_{k,R}} \|X^n - X^k\|_{\infty, T_0}^\alpha d\gamma_d \right) &= \int_{\Omega \times \mathbb{R}^d} \|X^n - X^k\|_{\infty, T_0}^\alpha d\hat{\Pi} \\ &= \int_{\Sigma_{n,k}^c} \|X^n - X^k\|_{\infty, T_0}^\alpha d\hat{\Pi} + \int_{\Sigma_{n,k}} \|X^n - X^k\|_{\infty, T_0}^\alpha d\hat{\Pi}. \end{aligned} \quad (5.22)$$

The first term on the right side of (5.22) is less than η^α , while the second one, due to (5.19) and (5.21), is dominated by

$$\sqrt{\hat{\Pi}(\Sigma_{n,k})} \cdot \sqrt{\mathbb{E} \int_{\mathbb{R}^d} \|X^n - X^k\|_{\infty, T_0}^{2\alpha} d\gamma_d} \leq 2^\alpha \sqrt{J_{\alpha, T_0} \hat{\Pi}(\Sigma_{n,k})} \rightarrow 0 \quad \text{as } n, k \rightarrow +\infty.$$

Now taking $\eta = \varepsilon^{1/\alpha}$ and combining (5.20) and (5.22), we prove that

$$\limsup_{n,k \rightarrow +\infty} \mathbb{E} \left[\int_{\mathbb{R}^d} \|X^n - X^k\|_{\infty, T_0}^\alpha d\gamma_d \right] \leq 2\varepsilon,$$

which implies (5.18).

Let $X \in L^\alpha(\Omega \times \mathbb{R}^d; C([0, T_0], \mathbb{R}^d))$ be the limit of X^n in this space. We see that for each $t \in [0, T]$ and almost all $x \in \mathbb{R}^d$, $w \rightarrow X_t(w, x)$ is in \mathcal{F}_t . \square

Proposition 5.4. *There exists a family $\{\hat{K}_t; t \in [0, T_0]\}$ of density functions on \mathbb{R}^d such that $(X_t)_\# \gamma_d = \hat{K}_t \gamma_d$ for each $t \in [0, T_0]$. Moreover, $\sup_{0 \leq t \leq T_0} \|\hat{K}_t\|_{L^p(\mathbb{P} \times \gamma_d)} \leq \Lambda_{p, T_0}$, where Λ_{p, T_0} is given in (5.16).*

Proof. It is the same as the proof of Theorem 3.4. \square

The same arguments in the proof of Proposition 4.1 and 4.2 yield the following

Proposition 5.5. *For any $\alpha \geq 2$, up to a subsequence,*

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^d} \mathbb{E} \left(\sup_{0 \leq t \leq T_0} \left| \sum_{i=1}^m \int_0^t [A_i^n(X_s^n) - A_i(X_s)] dw_s^i \right|^\alpha \right) d\gamma_d = 0,$$

and

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^d} \left[\mathbb{E} \int_0^{T_0} |A_0^n(X_s^n) - A_0(X_s)|^\alpha ds \right] d\gamma_d = 0.$$

Now for regularized vector fields $A_i^n, i = 0, 1, \dots, m$, we have

$$X_t^n(x) = x + \sum_{i=1}^m \int_0^t A_i^n(X_s^n) dw_s^i + \int_0^t A_0^n(X_s^n) ds. \quad (5.23)$$

When $n \rightarrow +\infty$, by Theorem 5.3 and Proposition 5.5, the two sides of (5.23) converge respectively to X and

$$x + \sum_{i=1}^m \int_0^t A_i(X_s) dw_s^i + \int_0^t A_0(X_s) ds$$

in the space $L^\alpha(\Omega \times \mathbb{R}^d; C([0, T_0], \mathbb{R}^d))$. Therefore for almost all $x \in \mathbb{R}^d$, the following equality holds \mathbb{P} -almost surely:

$$X_t(x) = x + \sum_{i=1}^m \int_0^t A_i(X_s) dw_s^i + \int_0^t A_0(X_s) ds, \quad \text{for all } t \in [0, T_0].$$

That is to say, X_t solves SDE (5.1) over $[0, T_0]$.

The following result proves the pathwise uniqueness to SDE (5.1) for a.e. initial value $x \in \mathbb{R}^d$.

Proposition 5.6. *Under the conditions (A1') and (A2)–(A4), the SDE (5.1) has a unique solution on the interval $[0, T_0]$.*

Proof. Let $(Y_t)_{t \in [0, T_0]}$ be another solution. Set, for $R > 0$,

$$G_R = \left\{ (w, x) \in \Omega \times \mathbb{R}^d; \sup_{0 \leq t \leq T_0} |X_t(w, x) - Y_t(w, x)| \leq R \right\}.$$

Remark that in Theorem 5.2, the terms involving $1/\sigma$ and $1/\sigma^2$ are equal to zero. Therefore the term

$$\begin{aligned} I &:= \mathbb{E} \int_{G_R} \log \left(\frac{\sup_{0 \leq t \leq T_0} |X_t - Y_t|^2}{\sigma^2} + 1 \right) d\gamma_d \\ &\leq C_{T_0} \Lambda_{p, T_0} C_{d, q, R} \left[\|A_0\|_{\mathbb{D}_1^q} + \left(\sum_{i=1}^m \|A_i\|_{\mathbb{D}_1^{2q}}^2 \right)^{\frac{1}{2}} + \sum_{i=1}^m \|A_i\|_{\mathbb{D}_1^{2q}}^2 \right] \end{aligned}$$

is bounded for any $\sigma > 0$. Consider for $\eta > 0$,

$$\Sigma_\eta = \left\{ (w, x); \sup_{0 \leq t \leq T_0} |X_t(w, x) - Y_t(w, x)| \geq \eta \right\}.$$

Similar to (5.21), we have

$$\mathbb{E} \left(\int_{G_R} \mathbf{1}_{\Sigma_\eta} d\gamma_d \right) \leq \frac{I}{\log(\frac{\eta^2}{\sigma^2} + 1)} \rightarrow 0$$

as $\sigma \rightarrow 0$. So we obtain

$$\mathbf{1}_{G_R} \cdot \sup_{0 \leq t \leq T_0} |X_t - Y_t| = 0, \quad (\mathbb{P} \times \gamma_d)\text{-a.s.}$$

Letting $R \rightarrow \infty$, we obtain that $(\mathbb{P} \times \gamma_d)$ almost surely, $X_t = Y_t$ for all $t \in [0, T_0]$. \square

Now we extend the solution to any time interval $[0, T]$. Let $\theta_{T_0} w$ be the time-shift of the Brownian motion w and denote by $X_t^{T_0}$ the corresponding solution to SDE driven by $\theta_{T_0} w$. By Proposition 5.6, $\{X_t^{T_0}(\theta_{T_0} w, x) : 0 \leq t \leq T_0\}$ is the unique solution to the SDE over $[0, T_0]$:

$$X_t^{T_0}(x) = x + \sum_{i=1}^m \int_0^t A_i(X_s^{T_0}(x)) d(\theta_{T_0} w)_s^i + \int_0^t A_0(X_s^{T_0}(x)) ds.$$

For $t \in [0, T_0]$, define $X_{t+T_0}(w, x) = X_t^{T_0}(\theta_{T_0} w, X_{T_0}(w, x))$. Note that X_t is well defined on the interval $[0, 2T_0]$ up to a $(\mathbb{P} \times \gamma_d)$ -negligible subset of $\Omega \times \mathbb{R}^d$. Replacing x by $X_{T_0}(x)$ in the above equation, we get easily

$$X_{t+T_0}(x) = x + \sum_{i=1}^m \int_0^{t+T_0} A_i(X_s(x)) dw_s^i + \int_0^{t+T_0} A_0(X_s(x)) ds.$$

Therefore X_t defined as above is a solution to SDE on the interval $[0, 2T_0]$. Continuing in this way, we obtain the solution of SDE (5.1) on $[0, T]$.

Theorem 5.7. *The $\{X_t; t \in [0, T]\}$ constructed above is the unique solution to SDE (5.1) in the sense of Definition 5.1. Moreover for each $t \in [0, T]$, the density K_t of $(X_t) \# \gamma_d$ with respect to γ_d is in the space $L^1 \log L^1$.*

Proof. Let Y_t , $t \in [0, T]$ be another solution in the sense of Definition 5.1. First by Proposition 5.6, we have $(\mathbb{P} \times \gamma_d)$ -almost surely, $Y_t = X_t$ for all $t \in [0, T_0]$. In particular, $Y_{T_0} = X_{T_0}$. Next by the flow property, Y_{t+T_0} satisfies the following equation:

$$Y_{t+T_0}(x) = Y_{T_0}(x) + \sum_{i=1}^m \int_0^t A_i(Y_{s+T_0}(x)) d(\theta_{T_0} w)_s^i + \int_0^t A_0(Y_{s+T_0}(x)) ds,$$

that is, Y_{t+T_0} is a solution with initial value Y_{T_0} . But by the above discussion, X_{t+T_0} is also a solution with the same initial value $X_{T_0} = Y_{T_0}$. Again by Proposition 5.6, we have $(\mathbb{P} \times \gamma_d)$ -almost surely, $X_{t+T_0} = Y_{t+T_0}$ for all $t \leq T_0$. Hence we have proved $X|_{[0, 2T_0]} = Y|_{[0, 2T_0]}$. Repeating this procedure, we obtain the uniqueness over $[0, T]$. The existence of density K_t of $(X_t) \# \gamma_d$ with respect to γ_d beyond T_0 is deduced from the flow property. However, to insure that $K_t \in L^1 \log L^1$, we have to use Theorem 3.3 and the following

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^d} \mathbb{E} \left(\sup_{0 \leq t \leq T} |X_t^n - X_t|^{\alpha} \right) d\gamma_d = 0,$$

which can be checked using the same arguments as in the proof of Propositions 4.1 and 4.2. \square

6 Appendix

For any locally integrable function $f \in L^1_{loc}(\mathbb{R}^d)$ and $R > 0$, the local maximal function $M_R f$ is defined by

$$M_R f(x) = \sup_{0 < r \leq R} \frac{1}{\text{Leb}_d(B(x, r))} \int_{B(x, r)} |f(y)| dy, \quad (6.1)$$

where $B(x, r) = \{y \in \mathbb{R}^d; |y - x| \leq r\}$. The following result is the starting point for the approach concerning Sobolev coefficients, used in [5] and [36].

Theorem 6.1. Let $f \in L^1_{loc}(\mathbb{R}^d)$ be such that $\nabla f \in L^1_{loc}(\mathbb{R}^d)$. Then there is a constant $C_d > 0$ (independent of f) and a negligible subset N , such that for $x, y \in N^c$ with $|x - y| \leq R$,

$$|f(x) - f(y)| \leq C_d |x - y| ((M_R |\nabla f|)(x) + (M_R |\nabla f|)(y)); \quad (6.2)$$

moreover for $p > 1$ and $f \in L^p_{loc}(\mathbb{R}^d)$, there is a constant $C_{d,p} > 0$ such that

$$\int_{B(r)} (M_R f)^p dx \leq C_{d,p} \int_{B(r+R)} |f|^p dx. \quad (6.3)$$

Since the inequality (6.2) played a key role in the proof of Theorem 5.2, we give here its proof for the sake of the reader's convenience.

We follow the idea of the proof of Claim #2 on p.253 in [9]. For any bounded measurable subset U in \mathbb{R}^d such that its Lebesgue measure $\text{Leb}_d(U) > 0$, define the average of $f \in L^1_{loc}(\mathbb{R}^d)$ on U by

$$(f)_U = \frac{1}{\text{Leb}_d(U)} \int_U f(y) dy.$$

Write $(f)_{x,r}$ instead of $(f)_{B(x,r)}$ for simplicity. Then $M_R f(x) = \sup_{0 < r \leq R} (f)_{x,r}$. We will need the following simple inequality: for any $C \in \mathbb{R}$,

$$|(f)_U - C| \leq \frac{1}{U} \int_U |f(y) - C| dy. \quad (6.4)$$

First, for any $x \in \mathbb{R}^d$ and $r \in]0, R]$, by Poincaré's inequality with $p = 1$ and $p^* = d/(d-1)$ (see [9] p.141), there is $C_d > 0$ such that

$$\begin{aligned} \int_{B(x,r)} |f - (f)_{x,r}| dy &\leq \left(\int_{B(x,r)} |f - (f)_{x,r}|^{d/(d-1)} dy \right)^{(d-1)/d} \\ &\leq C_d r \int_{B(x,r)} |\nabla f| dy \leq C_d M_R |\nabla f|(x) r. \end{aligned} \quad (6.5)$$

In particular, for all $k \geq 0$, by (6.4) and (6.5),

$$\begin{aligned} |(f)_{x,r/2^{k+1}} - (f)_{x,r/2^k}| &\leq \int_{B(x,r/2^{k+1})} |f - (f)_{x,r/2^k}| dy \\ &\leq 2^d \int_{B(x,r/2^k)} |f - (f)_{x,r/2^k}| dy \\ &\leq 2^d C_d M_R |\nabla f|(x) r/2^k. \end{aligned}$$

Since $f \in L^1_{loc}(\mathbb{R}^d)$, there is a negligible subset $N \subset \mathbb{R}^d$, such that for all $x \in N^c$, $f(x) = \lim_{r \rightarrow 0} (f)_{x,r}$. Thus for any $x \in N^c$, we have by summing up the above inequality that

$$|f(x) - (f)_{x,r}| \leq \sum_{k=0}^{\infty} |(f)_{x,r/2^{k+1}} - (f)_{x,r/2^k}| \leq 2^{1+d} C_d M_R |\nabla f|(x) r. \quad (6.6)$$

Next for all $x, y \in N^c, x \neq y$ and $|x - y| \leq R$, let $r = |x - y|$. Then by the triangular inequality, (6.4) and (6.5),

$$\begin{aligned} |(f)_{x,r} - (f)_{y,r}| &\leq \int_{B(x,r) \cap B(y,r)} (|(f)_{x,r} - f(z)| + |f(z) - (f)_{y,r}|) dz \\ &\leq \tilde{C}_d \left[\int_{B(x,r)} |(f)_{x,r} - f(z)| dz + \int_{B(y,r)} |f(z) - (f)_{y,r}| dz \right] \end{aligned}$$

$$\leq \tilde{C}_d C_d (M_R |\nabla f|(x) + M_R |\nabla f|(y)) r. \quad (6.7)$$

Now (6.2) follows from the triangular inequality and (6.6), (6.7):

$$\begin{aligned} |f(x) - f(y)| &\leq |f(x) - (f)_{x,r}| + |(f)_{x,r} - (f)_{y,r}| + |(f)_{y,r} - f(y)| \\ &\leq 2^{1+d} C_d M_R |\nabla f|(x) r + \tilde{C}_d C_d (M_R |\nabla f|(x) + M_R |\nabla f|(y)) r \\ &\quad + 2^{1+d} C_d M_R |\nabla f|(y) r \\ &= C_d (2^{1+d} + \tilde{C}_d) |x - y| (M_R |\nabla f|(x) + M_R |\nabla f|(y)). \end{aligned}$$

We obtain (6.2). \square

References

- [1] L. Ambrosio, Transport equation and Cauchy problem for BV vector fields. *Invent. Math.* 158 (2004), 227–260.
- [2] L. Ambrosio and A. Figalli, On flows associated to Sobolev vector fields in Wiener space: an approach à la DiPerna-Lions. *J. Funct. Anal.* 256 (2009), no. 1, 179–214.
- [3] L. Ambrosio, M. Lecumberry and S. Maniglia, Lipschitz regularity and approximate differentiability of the Di Perna-Lions flow. *Rend. Sem. Mat. Univ. Padova*, 114 (2005), 29–50.
- [4] F. Cipriano and A.B. Cruzeiro, Flows associated with irregular \mathbb{R}^d -vector fields. *J. Diff. Equations* 210 (2005), 183–201.
- [5] G. Crippa and C. De Lellis, Estimates and regularity results for the DiPerna-Lions flows. *J. Reine Angew. Math.* 616 (2008), 15–46.
- [6] A.B. Cruzeiro, Équations différentielles ordinaires: Non explosion et mesures quasi-invariantes. *J. Funct. Anal.* 54 (1983), 193–205.
- [7] R.J. Di Perna and P.L. Lions, Ordinary differential equations, transport theory and Sobolev spaces. *Invent. Math.* 98 (1989), 511–547.
- [8] B. Driver, Integration by parts and quasi-invariance for heat kernal measures on loop groups. *J. Funct. Anal.* 149 (1997), 470–547.
- [9] L.C. Evans and R.F. Gariepy, Measure theory and fine properties of functions. *Studies in Advanced Math.*, CRC Press, London, 1992.
- [10] S. Fang, Canonical Brownian motion on the diffeomorphism group of the circle. *J. Funct. Anal.* 196 (2002), 162–179.
- [11] S. Fang, P. Imkeller, T. Zhang, Global flows for stochastic differential equations without global Lipschitz conditions. *Ann. Probab.* 35 (2007), 180–205.
- [12] Shizan Fang and Dejun Luo, Transport equations and quasi-invariant flows on the Wiener space. *Bull. Sci. Math.* (2009), doi: 10.1016/j.bulsci.2009.01.001.
- [13] S. Fang, T. Zhang, A study of a class of stochastic differential equations with non-Lipschitzian coefficients. *Probab. Theory Related Fields* 132 (2005), 356–390.
- [14] A. Figalli, Existence and uniqueness of martingale solutions for SDEs with rough or degenerate coefficients. *J. Funct. Anal.* 254 (2008), 109–153.

- [15] F. Flandoli, M. Gubinelli and E. Priola, Well-posedness of the transport equation by stochastic perturbation. *Invent. Math.* (2009), doi: 10.1007/s00222-009-0224-4.
- [16] Zhiyuan Huang, Foundations of stochastic analysis (in Chinese). Second edition, Science Press of China, 2001.
- [17] N. Ikeda and S. Watanabe, Stochastic Differential Equations and Diffusion Processes. Second edition. North-Holland, Amsterdam, 1989.
- [18] H. Kaneko and S. Nakao, A note on approximation for stochastic differential equations. *Séminaire de Probabilités, XXII*, 155–162, Lecture Notes in Math., 1321, Springer, Berlin, 1988.
- [19] N.V. Krylov, On weak uniqueness for some diffusion with discontinuous coefficients. *Stochastic Process. Appl.* 113 (2004), 37–64.
- [20] N.V. Krylov and M. Röckner, Strong solutions of stochastic equations with singular time dependent drift. *Probab. Theory Related Fields* 131 (2005), 154–196.
- [21] H. Kunita, Stochastic Flows and Stochastic Differential Equations. Cambridge University Press, 1990.
- [22] C. LeBris and P.L. Lions, Existence and uniqueness of solutions to Fokker-Planck type equations with irregular coefficients. *Comm. Partial Differential Equations* 33 (2008), 1272–1317.
- [23] Y. Le Jan and O. Raimond, Integration of Brownian vector fields. *Ann. Probab.* 30 (2002), 826–873.
- [24] Y. Le Jan and O. Raimond, Flows, coalescence and noise. *Ann. Probab.* 32 (2004), 1247–1315.
- [25] X.M. Li, Strong p-completeness of stochastic differential equations and the existence of smooth flows on non-compact manifolds. *Probab. Theory Related Fields* 100 (1994), 485–511.
- [26] X.M. Li and M. Scheutzow, Lack of strong completeness for stochastic flows. <http://arxiv.org/abs/0908.1839>.
- [27] Dejun Luo, Quasi-invariance of Lebesgue measure under the homeomorphic flow generated by SDE with non-Lipschitz coefficient. *Bull. Sci. Math.* 133 (2009), 205–228.
- [28] P. Malliavin, Stochastic Analysis, *Grundl. der Math. Wissenschaften*, vol. 313, Springer, 1997.
- [29] P. Malliavin, The canonical diffusion above the diffeomorphism group of the circle. *C. R. Acad. Sci.* 329 (1999), 325–329.
- [30] D. Revuz and M. Yor, Continuous martingale and Brownian motion, *Grund. der Math. Wiss.* 293, 1991, Springer-Verlag.
- [31] M.V. Safonov, Non uniqueness for second order elliptic equations with measurable coefficients. *SIAM J. Math. Anal.* 30 (1999), 879–895.
- [32] D.W. Stroock and S.R.S Varadhan, Multidimensional diffusion processes, Springer, New York, 1979.

- [33] A.J. Veretennikov, On the strong solutions of stochastic differential equations. *Theory Prob. Appl.* 24 (1979), 354–366.
- [34] Xicheng Zhang, Strong solutions of SDEs with singular drift and Sobolev diffusion coefficients. *Stochastic Process. Appl.* 115 (2005), 1805–1818.
- [35] Xicheng Zhang, Homeomorphic flows for multi-dimensional SDEs with non-Lipschitz coefficients. *Stochastic Process. Appl.* 115 (2005), no. 3, 435–448; Erratum to “Homeomorphic flows for multi-dimensional SDEs with non-Lipschitz coefficients”. *Stochastic Process. Appl.* 116 (2006), no. 5, 873–875.
- [36] Xicheng Zhang, Stochastic flows of SDEs with irregular coefficients and stochastic transport equations. *Bull. Sci. Math.* (2009), doi:10.1016/j.bulsci.2009.12.004.